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05b. Keyhole/Hankel contour and $\zeta(-n)$

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The contour-integration trick illustrated here appeared in one of Riemann's proofs of analytic continuation of $\zeta(s)$. It almost immediately proves that values of $\zeta(s)$ at non-positive integers are *rational*, and shows the connection to the Laurent coefficients of $1/(e^t - 1)$ at t = 0.

[1.1] An integral representation of $\Gamma(s) \cdot \zeta(s)$ Although the integral representation of $\zeta(s)$ using a theta function is perhaps better in the long run, there is a more elementary one. As always,

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s} \qquad (\text{in } \operatorname{Re}(s) > 1)$$

[1.2] Claim: For $\operatorname{Re}(s) > 1$,

$$\Gamma(s) \cdot \zeta(s) = \int_0^\infty \frac{t^s}{e^t - 1} \frac{dt}{t}$$

Proof: Expand a geometric series, exchange sum and integral, and change variables:

$$\int_{0}^{\infty} \frac{t^{s}}{e^{t} - 1} \frac{dt}{t} = \int_{0}^{\infty} \frac{t^{s} e^{-t}}{1 - e^{-t}} \frac{dt}{t} = \int_{0}^{\infty} t^{s} \left(\sum_{n \ge 1} e^{-nt}\right) \frac{dt}{t} = \sum_{n \ge 1} \int_{0}^{\infty} t^{s} e^{-nt} \frac{dt}{t}$$
$$= \sum_{n \ge 1} \frac{1}{n^{s}} \int_{0}^{\infty} t^{s} e^{-t} \frac{dt}{t} = \Gamma(s) \cdot \sum_{n \ge 1} \frac{1}{n^{s}} = \Gamma(s) \cdot \zeta(s)$$
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as claimed.

[1.3] Keyhole/Hankel contour The *keyhole* or *Hankel* contour is a path from $+\infty$ inbound along the real line to $\varepsilon > 0$, counterclockwise around a circle of radius ε at 0, back to ε on the real line, and outbound back to $+\infty$ along the real line.

The usual elementary application is to evaluation of integrals similar to $\int_0^\infty \frac{t^s dt}{t^2+1}$, with $0 < \operatorname{Re}(s) < 1$. In such an example, analytically continuing counterclockwise around 0 has no impact on the denominator, but, significantly, the numerator changes by a factor $e^{2\pi i s}$, since

 $t^s = (|t| \cdot e^{i\theta})^s = |t|^s \cdot e^{i\theta s}$ (and θ goes from 0 to 2π)

We want the out-bound value of t^s to be real-valued for real s, so the inbound version of t^s must be actually be $t^s \cdot e^{2\pi i s}$.

The absolute value of the integrand goes to 0 as $|t| \to 0$, so the integral over the small circle goes to 0 as $\varepsilon \to 0$, as do the integrals to and from $0, \varepsilon$ along the real line.

Thus, letting H_{ε} be the Hankel contour with circle of radius $\varepsilon > 0$,

$$\lim_{\varepsilon \to 0} \int_{H_{\varepsilon}} \frac{t^s dt}{t^2 + 1} = \lim_{\varepsilon \to 0} \left(\int_{+\infty}^{\varepsilon} \frac{(t \cdot e^{2\pi i})^s dt}{t^2 + 1} + (\text{integral over little circle}) + \int_{\varepsilon}^{+\infty} \frac{t^s dt}{t^2 + 1} \right)$$
$$= (1 - e^{2\pi i s}) \int_{0}^{\infty} \frac{t^s dt}{t^2 + 1}$$

The standard way to make the Residue Theorem useful is to modify H_{ε} by not going all the way to $+\infty$ outbound, but stopping at +R for large positive R, traversing clockwise a large circle of radius R back to the positive real axis, and then inbound to ε . We anticipate that the integrals from R to and from $+\infty$ go to 0 as $R \to +\infty$, as does the integral over the large circle.

A touch of care is necessary to correctly estimate $z^s/(z^2+1)$ on |z|=R, with our choice that

$$(Re^{i\theta})^s = e^{s\log R} \cdot e^{is\theta} \qquad (\text{with } 0 \le \theta < 2\pi, R > 0, \log R \in \mathbb{R})$$

In particular, the reliable conventional fact that $|R^{it}| = 1$ for positive real R and real t, is inadequate for a treatment of exponentials of complex numbers. For our specification of what z^s means in the present context, with $z = Re^{i\theta}$, R > 0, $0 \le \theta < 2\pi$, and s = u + iv,

$$|z^{s}| = |(R \cdot e^{i\theta})^{s}| = |R^{s}| \cdot |e^{i\theta s}| = R^{\operatorname{Re}(s)} \cdot |e^{i\theta(u+iv)}|$$

Since $0 \le \theta < 2\pi$, $|e^{-\theta v}| \le e^{2\pi|v|}$. Since s = u + iv is fixed in this discussion, we have a uniform bound $C = e^{2\pi|v|}$. Thus, on the circle |z| = R,

$$|z^{s}| = |(R \cdot e^{i\theta})^{s}| \leq R^{\operatorname{Re}(s)} \cdot e^{2\pi |\operatorname{Im}(s)|}$$

Specifically, for fixed s with $-1 < \operatorname{Re}(s) < 1$, this is bounded by $C \cdot R^1$. Thus,

$$\left| \text{integral over big circle} \right| \le \ \text{length} \cdot \max \ \text{value} \le \ 2\pi R \cdot \frac{C \cdot R^{\text{Re}(s)}}{R^2 - 1}$$

For each R, ε , this gives a path integral (counter-clockwise) over a *closed* path. By *residues*, this picks up $2\pi i$ times the sum of the residues inside the path. Thus, we discover that the integrals do not depend on the parameters $0 < \varepsilon < 1 < R$. Keeping track of the relevant versions of t^s ,

$$(1 - e^{2\pi is}) \int_0^\infty \frac{t^s dt}{t^2 + 1} = 2\pi i \cdot \left((\text{residue at } t = i) + (\text{residue at } t = -i) \right)$$
$$= 2\pi i \cdot \left(+ \frac{e^{\frac{1}{2}\pi is}}{-i - i} + \frac{e^{\frac{3}{2}\pi is}}{i + i} \right) = \pi \cdot \left(e^{\frac{1}{2}\pi is} - e^{\frac{3}{2}\pi is} \right)$$

That is,

$$\int_0^\infty \frac{t^s \, dt}{t^2 + 1} = \pi \cdot \frac{e^{\frac{1}{2}\pi i s} - e^{\frac{3}{2}\pi i s}}{1 - e^{2\pi i s}} = \pi \cdot \frac{e^{-\frac{1}{2}\pi i s} - e^{\frac{1}{2}\pi i s}}{e^{-\pi i s} - e^{\pi i s}} = \frac{\pi}{e^{\frac{1}{2}\pi i s} + e^{-\frac{1}{2}\pi i s}} = \frac{2\pi}{\cos\frac{\pi s}{2}}$$

This is a charming and useful device, but a different secondary trick is applied to $\zeta(s)$:

[1.4] Evaluation of $\zeta(-n)$ The first part of the Hankel contour discussion gives

$$\Gamma(s) \cdot \zeta(s) = \int_0^\infty \frac{t^s}{e^t - 1} \frac{dt}{t} = \frac{1}{1 - e^{2\pi i (s-1)}} \cdot \lim_{\varepsilon \to 0} \int_{H_\varepsilon} \frac{t^s}{e^t - 1} \frac{dt}{t} = \frac{1}{1 - e^{2\pi i s}} \cdot \lim_{\varepsilon \to 0} \int_{H_\varepsilon} \frac{t^s}{e^t - 1} \frac{dt}{t}$$

Rewrite this as

$$\zeta(s) = \frac{1}{\Gamma(s) \cdot (1 - e^{2\pi i s})} \cdot \lim_{\varepsilon \to 0} \int_{H_{\varepsilon}} \frac{t^s}{e^t - 1} \frac{dt}{t}$$

At $s = -n \in \{0, -1, -2, -3, -4, ...\}$ two fortunate things happen. First, the pole of $\Gamma(s)$ and the zero of $1 - e^{2\pi i s}$ cancel, giving a finite, computable value. Second, the function t^{-n-1} is *single-valued*, so the inbound and outbound integrals of the Hankel contour simply *cancel* each other, *and* the integral over the small circle at 0 becomes $2\pi i$ times the residue of $\frac{t^{-n-1}}{e^{t-1}}$ at 0.

The periodicity of $1 - e^{2\pi i s}$ assures that the leading (linear) term in the power series at any integer is the same as that at 0, namely,

$$1 - e^{2\pi i s} = 1 - \left(1 + \frac{2\pi i s}{1!} + \frac{(2\pi i s)^2}{2!} + \dots\right) = 2\pi i s + \text{higher}$$

Grant for the moment that the residue of $\Gamma(s)$ at -n is $(-1)^n/n!$. Then

$$\zeta(-n) = \frac{1}{\frac{(-1)^n}{n!} \cdot 2\pi i} \cdot 2\pi i \cdot \operatorname{Res}_{t=0} \frac{t^{-n-1}}{e^t - 1} = (-1)^n \cdot n! \cdot \operatorname{Res}_{t=0} \frac{t^{-n-1}}{e^t - 1}$$

The Laurent coefficients of $\frac{t^{-n-1}}{e^t-1}$ are more-or-less Bernoulli numbers. These are not completely elementary objects, but are certainly *rational*. Thus, $\zeta(-n) \in \mathbb{Q}$.

[1.5] Vanishing $\zeta(-2) = \zeta(-4) = \ldots = 0$ A slightly finer analysis of the generating function $\frac{1}{e^t-1}$ yields the vanishing of $\zeta(s)$ at negative even integers, as follows.

First, $\frac{1}{e^t-1}$ is very close to being *odd* as a function of t:

$$\frac{1}{e^t - 1} + \frac{1}{e^{-t} - 1} = \frac{1}{e^t - 1} + \frac{e^t}{1 - e^t} = \frac{1}{e^t - 1} - \frac{e^t}{e^t - 1} = \frac{1 - e^t}{e^t - 1} = -1$$

Thus,

$$\left(\frac{1}{e^t - 1} + \frac{1}{2}\right) + \left(\frac{1}{e^{-t} - 1} + \frac{1}{2}\right) = 0$$

and $\frac{1}{e^t-1} + \frac{1}{2}$ is *odd*, so all its non-vanishing Laurent coefficients are odd-degree. Thus, for even -2n < 0,

$$\zeta(-2n) = (-1)^{2n} (2n)! \operatorname{Res}_{t=0} \frac{t^{-2n-1}}{e^t - 1} = (2n)! (2n^{th} \text{ Laurent coefficient of } \frac{1}{e^t - 1}) = 0$$

[1.6] Residues of $\Gamma(s)$ Finally, we determine the residues of $\Gamma(s)$. Certainly

$$\Gamma(1) = \int_0^\infty t^1 e^{-t} \frac{dt}{t} = \int_0^\infty e^{-t} dt = 1$$

From the functional equation $s\Gamma(s) = \Gamma(s+1)$, near s = 0

$$\Gamma(s) = \frac{\Gamma(s+1)}{s} = \frac{1 + \text{higher}}{s} = \frac{1}{s} + (\text{holomorphic at } s = 0)$$

Thus, the residue at 0 is 1. Iterating the functional equation,

$$\Gamma(s) = \frac{\Gamma(s+1)}{s} = \frac{\Gamma(s+2)}{(s+1)s} = \frac{\Gamma(s+3)}{(s+2)(s+1)s} = \dots = \frac{\Gamma(s+n+1)}{(s+n)(s+n-1)\dots(s+2)(s+1)s}$$

Thus, the leading Laurent term at s = -n is

$$\frac{1}{s+n} \cdot \frac{\Gamma(s+n+1)}{(s+n-1)\dots(s+2)(s+1)s}\Big|_{s=-n} = \frac{1}{s+n} \cdot \frac{\Gamma(-n+n+1)}{(-n+n-1)\dots(-n+2)(-n+1)(-n)}$$
$$= \frac{1}{s+n} \cdot \frac{1}{(-1)(-2)(-3)\dots(-n+2)(-n+1)(-n)} = \frac{1}{s+n} \cdot \frac{(-1)^n}{n!}$$

That is, the residue of $\Gamma(s)$ at -n is $(-1)^n/n!$ as claimed.

Bibliography

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