

(November 1, 2012)

Functional analysis exercises-discussion 02

Paul Garrett garrett@math.umn.edu <http://www.math.umn.edu/~garrett/>

[This document is http://www.math.umn.edu/~garrett/m/fun/exercises_2012-13/fun-disc-10-27-2012.pdf]

Was due Wed, 24 Oct 2012, preferably as PDF emailed to me.

[02.1] *Convincingly* and *not-ugly-ly* prove that e^{-1/x^2} (naturally extended by 0 at $x = 0$) is infinitely differentiable at 0.

Using $e^{-a} = 1/e^a$,

$$e^{-1/x^2} = \frac{1}{e^{1/x^2}} = \frac{1}{1 + \frac{1}{x^2} + \dots + \frac{1}{\ell! \cdot x^{2\ell}} + \dots} \leq \frac{1}{\frac{1}{\ell! \cdot x^{2\ell}}} = \ell! \cdot x^{2\ell} \quad (\text{for every } 1 \leq \ell \in \mathbb{Z})$$

A trivial induction step shows that (at $x \neq 0$) $\frac{d^n}{dx^n} e^{-1/x^2}$ is of the form $R(x) \cdot e^{-1/x^2}$ for some rational function R depending on n . Thus, with these derivatives inductively extending to $x = 0$ by taking value 0 there,

$$\left| \frac{R(x)e^{-1/x^2} - 0}{x} \right| \leq \left| \frac{R(x) \cdot \ell! \cdot x^{2\ell}}{x} \right| \rightarrow 0 \quad (\text{as } x \rightarrow 0, \text{ for } \ell \text{ large enough})$$

Thus, the $(n+1)$ st derivative is also 0 at 0, and the induction succeeds. ///

[02.2] Show that for $0 \leq c_n \in \mathbb{R}$ with c_n decreasing *monotonically* to 0, the Fourier series $\sum_n c_n e^{inx}$ converges at $x \notin 2\pi\mathbb{Z}$, although not necessarily *absolutely*.

Yes, this is a special case of *summation by parts*, since all we really need is that the partial sums $\sum_{|n| \leq N} e^{inx}$ be *bounded*, which holds for fixed $x \notin 2\pi\mathbb{Z}$, since that finite sum a geometric series, which sums to $(e^{-inx} - e^{i(n+1)x})/(1 - e^{ix})$.

Suppose, generally, that the finite subsums $B_N = \sum_{|n| \leq N} b_n$ are *bounded*, and c_n 's positive real, decreasing monotonically to 0. Let's break the sum into two pieces, one with $n \geq 0$ and the other with $n < 0$, to avoid clumsiness of language. We estimate the tails of the sum: partial summation gives

$$\sum_{M \leq n \leq N} b_n c_n = b_M c_M + \sum_{M+1 \leq \ell \leq N} (B_\ell - B_{\ell-1}) \cdot c_\ell = b_M c_M + \sum_{M+1 \leq \ell \leq N-1} B_\ell \cdot (c_\ell - c_{\ell+1}) + B_N c_N$$

The first and last terms $b_M c_M$ and $B_N c_N$ go to 0, because b_M and B_N are bounded, and c_M and c_N go to 0. Letting B be a finite bound for the B_N 's, the middle sum is estimated by

$$\left| \sum_{M+1 \leq \ell \leq N-1} B_\ell \cdot (c_\ell - c_{\ell+1}) \right| \leq B \cdot \sum_{M+1 \leq \ell \leq N-1} |c_\ell - c_{\ell+1}| = B \cdot \sum_{M+1 \leq \ell \leq N-1} c_\ell - c_{\ell+1} = B \cdot (c_{M+1} - c_N)$$

using the monotonicity of the c_ℓ 's. This, too, goes to 0. That is, the finite tails go to 0, so the infinite sum converges. ///

[02.3] The Fejér kernel discussion was carried out in considerable detail in class, I think.

[02.4] Map $\mathbb{C}^n \rightarrow \mathbb{C}^{n+1}$ by $(z_1, \dots, z_n) \rightarrow (z_1, \dots, z_n, 0)$. Let $V = \text{colim}_n \mathbb{C}^n = \bigcup_n \mathbb{C}^n$, with the colimit topology, in which a basis of opens at 0 is given by convex hulls of unions $B = \bigcup_n B_n$ where B_n is an open

ball of some positive radius, at 0, in \mathbb{C}^n . Show that V violates the conclusion of the Baire category theorem, so is not complete-metrizable. Here *Cauchy sequences* $\{x_n\}$ are those such that, given a neighborhood N of 0, there is n_o such that $x_m - x_n \in N$ for all $m, n \geq n_o$. Show that Cauchy sequences *converge*.

Each subspace \mathbb{C}^n is nowhere dense in the colimit, because elements of \mathbb{C}^n cannot come close to elements of \mathbb{C}^N with $N > n$ and non-zero components at $(n+1)$ st and subsequent entries. Thus, $\text{colim}_n \mathbb{C}^n$ is a countable union of nowhere dense subsets, violating the conclusion of the Baire category theorem. Thus, this colimit is not complete-metrizable.

We claim that in the colimit (*diamond*) topology, every Cauchy sequence in fact lies in some limitand \mathbb{C}^n , the latter with its usual complete topology, so the sequence converges.

Indeed, let v_1, v_2, \dots be a sequence *not* lying in any single limitand \mathbb{C}^n . Without loss of generality, replace the sequence by a subsequence so that $v_\ell \in \mathbb{C}^{n_\ell}$ but *not* in $\mathbb{C}^{n_\ell-1}$, with a strictly increasing sequence $n_1 < n_2 < n_3 < \dots$. Let

$$v_\ell = (v_{\ell 1}, v_{\ell 2}, \dots, v_{\ell n_\ell}, 0, 0, \dots)$$

Note that $v_{\ell n_\ell} \neq 0$. Let

$$B_\ell = \text{open ball of radius } \min_{j \leq \ell} \frac{|v_{jn_j}|}{2} \subset \mathbb{C}^{n_\ell}$$

Then $v_\ell - v_{\ell-1}$ has n_ℓ^{th} component $v_{\ell n_\ell}$, so is *not* in B_ℓ . Let B be the convex hull of the union of the B_ℓ s. We claim that $v_\ell - v_{\ell-1}$ is *not* in B , for all ℓ . This would prove that the sequence v_ℓ is not Cauchy.

Since $B_{\ell'} \cap \mathbb{C}^{n_\ell} \subset B_\ell$, $v_\ell - v_{\ell-1}$ is not in $B_{\ell'}$ for any $\ell' \geq \ell$. And taking the convex hull with $B_{\ell'}$ with $\ell' < \ell$ cannot affect the n_ℓ^{th} entries, since $B_{\ell'} \subset \mathbb{C}^{n_\ell-1}$ for $\ell' < \ell$. Thus, for every ℓ , $v_\ell - v_{\ell-1}$ is not in the neighborhood B of 0, so the sequence v_ℓ is not Cauchy. ///

[02.5] Show that for complex w the equation $(\frac{d^2}{dx^2} - w^2) u_w = 0$ has a solution in $C^2(S^1)$ only when $w \in i\mathbb{Z}$. Let δ^{per} be the $2\pi\mathbb{Z}$ periodic Dirac δ -function, the continuous linear functional on $C^o(\mathbb{R}/2\pi\mathbb{Z})$ given by $\delta^{\text{per}} f = f(0)$. With $w \in \mathbb{C}$ and $w \notin i\mathbb{Z}$, solve

$$\left(\frac{d^2}{dx^2} - w^2\right) u_w = \delta^{\text{per}}$$

for u_w on $\mathbb{R}/2\pi\mathbb{Z}$. (*Hint*: expand δ^{per} in a Fourier series.) Identify the residues of the $L^2(S^1)$ -valued meromorphic function $w \rightarrow u_w$.

The differential equation $(\frac{d^2}{dx^2} - w^2) u_w = 0$ on \mathbb{R} has solutions $u_w(x) = Ae^{wx} + Be^{-wx}$ for arbitrary scalars A, B , for $w \neq 0$, and $A + Bx$ for $w = 0$. (The fact that there are no *other* solutions follows from the Mean Value Theorem, for example.) Such a solution *descends* to the quotient $S^1 = \mathbb{R}/2\pi\mathbb{Z}$ only when it is invariant under translation by $2\pi\mathbb{Z}$, which requires $w \in i\mathbb{Z}$, and when $w = 0$ the solution $u(x) = x$ is excluded.

The Fourier series for periodic δ converges in $H^{-\frac{1}{2}-\varepsilon}(S^1)$ for every $\varepsilon > 0$. The equation

$$\sum_n 1 \cdot e^{inx} = \delta^{\text{per}} = \left(\frac{d^2}{dx^2} - w^2\right) \left(\sum_n \hat{u}(n) \cdot e^{inx}\right) = \sum_n (-n^2 - w^2) \cdot \hat{u}(n) \cdot e^{inx} \quad (\text{in } H^{-1}(S^1))$$

with differentiation extended by continuity to a continuous linear map on Levi-Sobolev spaces, gives $\hat{u}(n) = -1/(n^2 + w^2)$. For w bounded away from $i\mathbb{Z}$ the series

$$u_w = -\sum_n \frac{1}{n^2 + w^2} \cdot e^{inx}$$

is convergent in $H^{\frac{3}{2}-\varepsilon}(S^1)$ for all $\varepsilon > 0$, uniformly on compacts. We claim that this is a complex-differentiable $H^{\frac{3}{2}-\varepsilon}(S^1)$ -valued function. Estimate the $H^{\frac{3}{2}-\varepsilon}(S^1)$ -norm of the difference between the natural difference quotient and the natural expression for the derivative:

$$\begin{aligned} \left| \frac{u_{w+h}(x) - u_w(x)}{h} - \sum_n \frac{-2w}{(n^2 + w^2)^2} \cdot e^{inx} \right|_{H^{\frac{3}{2}-\varepsilon}(S^1)}^2 &\leq \sum_n \left| \frac{\frac{1}{n^2+(w+h)^2} - \frac{1}{n^2+w^2}}{h} + \frac{2w}{(n^2 + w^2)^2} \right|^2 \cdot (1+n^2)^{\frac{3}{2}-\varepsilon} \\ &= \sum_n \left| \frac{1}{h(n^2 + (w+h)^2)} - \frac{1}{h(n^2 + w^2)} + \frac{2w}{(n^2 + w^2)^2} \right|^2 \cdot (1+n^2)^{\frac{3}{2}-\varepsilon} \\ &= \sum_n \left| \frac{-2w-h}{(n^2 + (w+h)^2)(n^2 + w^2)} + \frac{2w}{(n^2 + w^2)^2} \right|^2 \cdot (1+n^2)^{\frac{3}{2}-\varepsilon} \\ &= h \cdot \sum_n \left| \frac{n^2 - 3w^2 - 2wh}{(n^2 + (w+h)^2)(n^2 + w^2)^2} \right|^2 \cdot (1+n^2)^{\frac{3}{2}-\varepsilon} \rightarrow 0 \end{aligned}$$

as $h \rightarrow 0$, uniformly for w in compacts away from $i\mathbb{Z}$, proving the complex differentiability. Thus, the only possible poles are at $w \in i\mathbb{Z}$, and are simple except for the double pole at $w = 0$.

Anticipating, as we will soon prove, that the elementary Cauchy-Goursat complex function theory applies to vector-valued functions: at $0 \neq in \in i\mathbb{Z}$ there are exactly two terms in the Fourier series with a pole, namely $e^{\pm inx}/(n^2 + w^2)$, with residues $e^{\pm inx}/(\pm 2in)$. At 0, the leading term of $1/w^2$ gives the constant 1, and the residue is 0.

Charmingly, apart from the double pole at $w = 0$, the residues of u_w produce eigenfunctions! ///

[02.6] Show that there are *no eigenvectors* for the Volterra operator $T : L^2[0, 1] \rightarrow L^2[0, 1]$ given by

$$Tf(x) = \int_0^x f(y) dy$$

By design, $\frac{d}{dx}Tf = f$ for $f \in C^o[0, 1]$. Show that $(T - \lambda)u = f$ is solvable for u when $0 \neq \lambda \in \mathbb{C}$ and $f \in C^1[0, 1]$ with $f(0) = 0$, and the solution is unique. With T^* the Hilbert-space adjoint of T , show that

$$(TT^*f)(x) = \int_0^1 \min(x, y) \cdot f(y) dy \quad (T^*Tf)(x) = \int_0^1 (1 - \max(x, y)) \cdot f(y) dy$$

Find eigenvectors for TT^* and T^*T .

First, a relation $Tf = \lambda \cdot f$ for $f \in L^2$ and $\lambda \neq 0$ would imply f is *continuous*:

$$|\lambda| \cdot |f(x+h) - f(x)| = |Tf(x+h) - Tf(x)| \leq \int_x^{x+h} 1 \cdot |f(t)| dt \leq |h|^{\frac{1}{2}} \cdot |f|_{L^2}$$

Then the fundamental theorem of calculus would imply f is continuously differentiable and $\lambda \cdot f' = (Tf)' = f$. Thus, f would be a constant multiple of $e^{x/\lambda}$, by the mean value theorem. However, again from Cauchy-Schwarz-Bunyakovsky, for a λ -eigenfunction

$$|\lambda| \cdot |f(x)| \leq |x|^{\frac{1}{2}} \cdot |f|_{L^2}$$

No non-zero multiple of these exponentials satisfies this equation. Thus, there are no eigenvectors for non-zero eigenvalues.

For $f \in L^2[0, 1]$ and $Tf = 0 \in L^2[0, 1]$, Tf is almost everywhere 0. In fact, since $x \rightarrow Tg(x)$ is unavoidably continuous, $Tf(x)$ is 0 for all x . Thus, for all x, y in the interval,

$$0 = 0 - 0 = Tf(y) - Tf(x) = \int_x^y f(t) dt$$

That is, $x \rightarrow Tf(x)$ is orthogonal in $L^2[0, 1]$ to all characteristic functions of intervals. Finite linear combinations of these are dense in $C^0[0, 1]$, which is dense in $L^2[0, 1]$. Thus $f = 0$. There are no eigenvectors for the Volterra operator.

One way to solve $(T - \lambda)u = f$ uses the idea that under some circumstances it makes sense to write

$$u = -\frac{1}{\lambda} \left(1 - \frac{T}{\lambda}\right)^{-1} f = -\frac{1}{\lambda} \left(1 + \frac{T}{\lambda} + \frac{T^2}{\lambda^2} + \dots\right) f$$

expanding the inverse as a geometric series. Convergence of the infinite sum will follow from estimating the operator norm of T^n . Indeed,

$$\begin{aligned} T^n f(x) &= \int_0^x \int_0^{x_{n-1}} \dots \int_0^{x_2} \int_0^{x_1} f(t) dt dx_1 \dots dx_{n-1} = \int_0^x f(t) \left(\int_t^x \int_t^{x_{n-1}} \dots \int_t^{x_2} dx_1 \dots dx_{n-1} \right) dt \\ &= \int_0^x f(t) \cdot \frac{(x-t)^{n-1}}{(n-1)!} dt \end{aligned}$$

From this, the operator norm of T^n is at most $1/n!$, so the geometric series expressing $(1 - T/\lambda)^{-1}$ converges for all $\lambda \neq 0$. The expression for T^n also gives

$$\left(1 + \frac{T}{\lambda} + \frac{T^2}{\lambda^2} + \dots\right) f(x) = \int_0^x f(t) \cdot \left(1 + \sum_{n \geq 1} \left(\frac{x-t}{\lambda}\right)^{n-1} / (n-1)!\right) dt = f(x) + \int_0^x f(t) \cdot e^{(x-t)/\lambda} dt$$

Thus, for $\lambda \neq 0$, we can solve $(T - \lambda)u = f$ by

$$u(x) = -\frac{f(x)}{\lambda} \int_0^x f(t) \cdot e^{(x-t)/\lambda} dt$$

Absence of eigenvectors assures *uniqueness*. For $\lambda = 0$, we need $f(0) = 0$, and then $u = f'$.

The compositions TT^* and T^*T are readily expressed by rewriting the operator T and its adjoint T^* as integrals against *kernel functions* $K(x, y)$, as

$$Tf(x) = \int_0^1 K(x, y) f(y) dy \quad \left(\text{where } K(x, y) = \begin{cases} y & (\text{for } y < x) \\ 0 & (\text{for } y > x) \end{cases}\right)$$

The kernel $K^*(x, y)$ of the adjoint is simply $K^*(x, y) = \overline{K(y, x)}$, as a direct computation shows:

$$\int_0^1 \left(\int_0^1 K(x, y) f(y) dy \right) \bar{g}(x) dx = \int_0^1 \int_0^1 K(x, y) f(y) \bar{g}(x) dx dy = \int_0^1 f(y) \overline{\int_0^1 K(x, y) g(x) dx} dy$$

Immediately

$$Tf(x) = \int_0^1 K(x, y) f(y) dy \quad \left(\text{where } K(x, y) = \begin{cases} 1 & (\text{for } y < x) \\ 0 & (\text{for } y > x) \end{cases}\right)$$

Then

$$K^*(x, y) = \overline{K(y, x)} = \begin{cases} 1 & (\text{for } x < y) \\ 0 & (\text{for } x > y) \end{cases} = \begin{cases} 1 & (\text{for } y > x) \\ 0 & (\text{for } y < x) \end{cases}$$

Thus,

$$\begin{aligned} TT^*f(x) &= \int_0^1 K(x,y) \int_0^1 K^*(y,t) f(t) dt dy = \int_0^x \int_y^1 f(t) dt dy \\ &= \int_0^1 f(t) \left(\int_0^{\min(x,t)} dy \right) dt = \int_0^1 \min(x,t) \cdot f(t) dt \end{aligned}$$

A similar computation works for T^*T .

As with T itself, $TT^*f = \lambda \cdot f$ with $\lambda \neq 0$ implies f is C^∞ , and $f = -\lambda \cdot f''$, so f is a linear combination of $e^{\pm\xi x}$ where $-\xi^2 = -1/\lambda$. Computing directly:

$$\begin{aligned} TT^*e^{\xi x} &= \int_0^x y \cdot f(y) dy + x \cdot \int_x^1 f(y) dy = \left[y \cdot \frac{e^{\xi y}}{\xi} \right]_0^x - \int_0^x \frac{e^{\xi y}}{\xi} dy + x \cdot \frac{e^{\xi \cdot 1} - e^{\xi x}}{\xi} \\ &= -\frac{e^{\xi x}}{\xi^2} + \frac{1}{\xi^2} + x \cdot \frac{e^{\xi \cdot 1}}{\xi} \end{aligned}$$

The second and third terms are the issue: a linear combination $Ae^{\xi x} + Be^{-\xi x}$ is a $1/\lambda$ -eigenvalue if and only if

$$A \cdot \left(\frac{1}{\xi^2} + x \cdot \frac{e^{\xi \cdot 1}}{\xi} \right) + B \cdot \left(\frac{1}{(-\xi)^2} + x \cdot \frac{e^{-\xi \cdot 1}}{-\xi} \right) = 0$$

which is

$$A + B = 0 \quad \text{and} \quad Ae^{\xi} - Be^{-\xi} = 0$$

Thus, $B = -A$, and then $e^{\xi} + e^{-\xi} = 0$, so $\xi \in \pi i\mathbb{Z} + \frac{\pi i}{2}$. That is, for $n \in \mathbb{Z}$, the unique (up to constant multiplies) $(\pi in + \frac{\pi i}{2})^2$ -eigenvector for TT^* is

$$u(x) = \sin \pi(n + \frac{1}{2})x \quad (\text{eigenvector for } TT^*)$$

The analogous computation can be done for T^*T , but, alternatively, we can observe that applying T^* to $TT^*f = \lambda \cdot f$ gives $(T^*T)(Tf) = \lambda \cdot (Tf)$. That is, eigenvectors for TT^* are mapped by T to eigenvectors for T^*T . Either way,

$$u(x) = \cos \pi(n + \frac{1}{2})x \quad (\text{eigenvector for } T^*T)$$

///

[0.0.1] Remark: Since T is *Hilbert-Schmidt*, it is provably *compact*, but $TT^* \neq T^*T$. Thus, we cannot expect T to have eigenvectors, and, indeed, it has none. Nevertheless, TT^* and T^*T *do* have eigenvectors. The same differential equation $f = -\lambda f''$ arises from both $TT^*f = \lambda \cdot f$ and $T^*Tf = \lambda \cdot f$, but, in effect, the implicit boundary conditions differ, so the eigenvectors differ.