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Introduction to zeta integrals and L -functions for GL_n

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1. Fourier-Whittaker expansions of cuspforms on GL_r
2. The Hecke-type case $GL_n \times GL_{n-1}$
3. The Rankin-Selberg case $GL_n \times GL_n$
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We give a quick introduction to Fourier-Whittaker expansions of cuspforms on GL_n , and integral representations of associated L -functions, following a part of Jacquet, Piatetski-Shapiro, and Shalika's extensions of Hecke's work for GL_2 .

All known ways to analytically continue automorphic L -functions involve *integral representations* using the corresponding automorphic forms. The simplest cases, extending Hecke's treatment of GL_2 , need no further analytic devices and very little manipulation beyond Fourier-Whittaker expansions. [1] Poisson summation is a sufficient device for several accessible classes of examples, as in Riemann, [Hecke 1918,20], [Tate 1950], [Iwasawa 1952], and [Godement-Jacquet 1972], and including treatment of the degenerate Eisenstein series needed for the $GL_n \times GL_n$ Rankin-Selberg convolutions. [2]

For f a cuspform on GL_n the most natural L -function obtained by an integral representation is the *Hecke-type* integral representation, also involving a cuspform F on GL_{n-1} ,

$$\Lambda(s, f \otimes F) = \int_{GL_{n-1}(k) \backslash GL_{n-1}(\mathbb{A})} |\det h|^{s-\frac{1}{2}} \cdot f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix} \cdot F(h) dh$$

More properly, the integral is a *zeta integral* $Z(s, f \times F)$, since within a given automorphic representation there is an essentially unique choice of automorphic form giving the correct local factors everywhere locally in that zeta integral: see [Jacquet-PS-Shalika 1981], [Jacquet-Shalika 1990], [Cogdell-PS 2003]. At *good* primes this is not an issue, but the general case must subsume the theory of *newforms*, as well as coping with complications at archimedean places. [3]

Another simple natural case, $m = n$, is the *Rankin-Selberg* integral using an auxiliary (degenerate) Eisenstein series

$$E_s(g) = \sum_{\gamma \in P_k^{n-1,1} \backslash GL_n(k)} \varphi_s(\gamma \cdot g)$$

where $P^{n-1,1}$ has Levi component $GL_{n-1} \times GL_1$, and

$$\varphi_s = \bigotimes_v \varphi_{s,v}$$

with $\varphi_{s,v}$ in a (degenerate) induced representation from $P_v^{n-1,1}$. The zeta integral attached to two cuspforms f, F on GL_n is

$$\Lambda(s, f \otimes F) = \int_{GL_n(k) \backslash GL_n(\mathbb{A})} E_s(h) \cdot f(h) \cdot F(h) dh$$

[1] See [Hecke 1937a,b]. Apparently the extension to $GL_{n-1} \times GL_n$ was considered so apparent that it was not explicitly mentioned in [Jacquet-PS-Shalika 1979], which was concerned with $GL_1 \times GL_3$ as prototype for $GL_m \times GL_n$.

[2] We do not discuss examples relying on meromorphic continuation of non-trivial Eisenstein series, as in [Langlands 1967/1976, 1971] and [Shahidi 1978,1985] have other requirements. See [Shahidi 2010] for a recent survey.

[3] In contrast to [Godement-Jacquet 1972], the *standard* L -function for f is *not* produced by this Hecke-type integral representation, except for $n = 2$. This was understood by Jacquet, Piatetski-Shapiro, and Shalika in the late 1970's, who developed the desired integral representations of a family of L -functions including the standard ones in the papers in the bibliography below. See also Cogdell's lecture notes in the bibliography.

1. Fourier-Whittaker expansions of cuspforms on GL_r

Some non-trivial aspects of the group structure of GL_r enters in the derivation of the Fourier expansion. The outcome is not obvious for $r > 2$.

Let $G = GL_r$, reserving the character n for elements of unipotent subgroups. Let

$$N^{\min} = \begin{pmatrix} 1 & * & \dots & * \\ & 1 & & \vdots \\ & & \ddots & * \\ 0 & & & 1 \end{pmatrix} = \text{unipotent radical of standard minimal parabolic}$$

Fix a non-trivial additive character ψ_o on $k \setminus \mathbb{A}$, and let ψ_{std} be the corresponding standard character on the unipotent radical of the standard minimal parabolic, namely,

$$\psi_{\text{std}}(u) = \psi_o(\text{sum super-diagonal entries}) = \psi_o(u_{12} + u_{23} + \dots + u_{r-1,r}) \quad (\text{non-trivial } \psi_o \text{ on } k \setminus \mathbb{A})$$

We obtain the Fourier expansion of a *cusppform* by an induction. First, a cusppform f has a Fourier expansion along the *abelian* unipotent radical

$$N = N^{r-1,1} = \{n_x = \begin{pmatrix} 1_{r-1} & x \\ 0 & 1 \end{pmatrix} : x = (r-1)\text{-by-1}\}$$

of the form

$$f(g) = \sum_{\psi} \int_{N_k \setminus N_{\mathbb{A}}} \bar{\psi}(n) f(ng) \, dn \quad (\psi \text{ summed over characters on } N_k \setminus N_{\mathbb{A}})$$

The cusppform condition implies that the component for the *trivial* character on $N_k \setminus N_{\mathbb{A}}$ is 0. The fragment

$$H = H^{r-1} = \left\{ \begin{pmatrix} GL_{r-1} & 0 \\ 0 & 1 \end{pmatrix} \right\}$$

of the Levi component of the parabolic $P^{r-1,1}$ acts *transitively* on the *non-trivial* characters on $N_k \setminus N_{\mathbb{A}}$. Letting

$$\psi_1(n_x) = \psi_o(x_{r-1})$$

the isotropy subgroup $\Theta = \Theta^{r-1}$ of ψ_1 in H is

$$\Theta = \{m \in H_k : \psi_1(mnm^{-1}) = \psi_1(n) \text{ for all } n \in N_{\mathbb{A}}\} = \left\{ \begin{pmatrix} GL_{r-2} & * & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \right\}$$

Thus, for a cusppform f ,

$$f(g) = \sum_{\gamma \in \Theta_k \setminus H_k} \int_{N_k \setminus N_{\mathbb{A}}} \bar{\psi}_1(\gamma n \gamma^{-1}) f(ng) \, dn$$

Replacing n by $\gamma^{-1}n\gamma$ and using the left G_k -invariance of f , this is

$$f(g) = \sum_{\gamma \in \Theta_k \setminus H_k} \int_{N_k \setminus N_{\mathbb{A}}} \bar{\psi}_1(n) f(n\gamma g) \, dn$$

For the induction step, let

$$N' = \{u_x = \begin{pmatrix} 1_{r-2} & x & 0 \\ 0 & 1 & 0 \\ 0 & & 1 \end{pmatrix}\} \subset H^{r-1}$$

Note that N' normalizes N , and

$$\psi_1(unu^{-1}) = \psi_1(n) \quad (\text{for all } n \in N_{\mathbb{A}} \text{ and } u \in N'_{\mathbb{A}})$$

Letting

$$H^{r-2} = \left\{ \begin{pmatrix} GL_{r-2} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \right\}$$

we have $\Theta = N'H^{r-2} = H^{r-2}N'$. For each γ , the function

$$h \longrightarrow \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nh\gamma g) dn \quad (\text{for } h \in H^{r-1})$$

is left N'_k -invariant, because

$$\int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(n\alpha h\gamma g) dn = \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(\alpha n h\gamma g) dn = \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nh\gamma g) dn \quad (\text{for } \alpha \in U_k)$$

by replacing n by $\alpha n \alpha^{-1}$ and using the left G_k -invariance of f . Thus, for each γ , there is a Fourier expansion along N' , namely,

$$\int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nh\gamma g) dn = \sum_{\psi'} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'(u) \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nuh\gamma g) dn du \quad (\text{characters } \psi' \text{ of } N'_k \backslash N'_{\mathbb{A}})$$

In fact, we only need $h = 1$:

$$\int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(n\gamma g) dn = \sum_{\psi'} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'(u) \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nu\gamma g) dn du \quad (\text{characters } \psi' \text{ of } N'_k \backslash N'_{\mathbb{A}})$$

The $\psi' = 1$ summand is 0, because f is cuspidal, since

$$N' \cdot (\ker \psi_1 \text{ on } N) \supset \left\{ \begin{pmatrix} 1_{r-2} & * & * \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \right\} = \text{unipotent radical of } (r-2, 2) \text{ parabolic}$$

The action of H^{r-2} on non-trivial characters ψ' on N' is *transitive*. Let

$$\psi'_1 \begin{pmatrix} 1_{r-2} & x & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \psi'_o(x_{r-2}) \quad (\text{with } x \text{ } (r-2)\text{-by-1})$$

The isotropy group of ψ'_1 is

$$\Theta^{r-2} = \left\{ \begin{pmatrix} GL_{r-3} & * & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \right\}$$

Thus,

$$\begin{aligned} \sum_{\psi'} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'(u) \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nu\gamma g) dn du &= \sum_{\delta \in \Theta_k^{r-2} \backslash H_k^{r-2}} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'_1(\delta u \delta^{-1}) \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(nu\gamma g) dn du \\ &= \sum_{\delta \in \Theta_k^{r-2} \backslash H_k^{r-2}} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'_1(u) \int_{N_k \backslash N_{\mathbb{A}}} \bar{\psi}_1(n) f(n\delta^{-1}u\delta\gamma g) dn du \end{aligned}$$

by replacing u by $\delta^{-1}u\delta$. We can also replace n by $\delta^{-1}n\delta$ without affecting ψ_1 , so this becomes

$$\sum_{\delta \in \Theta_k^{r-2} \backslash H_k^{r-2}} \int_{N'_k \backslash N'_\mathbb{A}} \bar{\psi}'_1(u) \int_{N_k \backslash N_\mathbb{A}} \bar{\psi}_1(n) f(nu\delta\gamma g) dn du$$

Altogether,

$$f(g) = \sum_{\gamma \in \Theta_k^{r-1} \backslash H_k^{r-1}} \sum_{\delta \in \Theta_k^{r-2} \backslash H_k^{r-2}} \int_{N'_k \backslash N'_\mathbb{A}} \bar{\psi}'_1(u) \int_{N_k \backslash N_\mathbb{A}} \bar{\psi}_1(n) f(nu\delta\gamma g) dn du$$

Since $\Theta_k^{r-1} = H_k^{r-2}N'$, the elements $\delta\gamma$ with $\gamma \in \Theta_k^{r-1} \backslash H_k^{r-1}$ and $\delta \in \Theta_k^{r-2} \backslash H_k^{r-2}$ are in natural bijection with $\Theta_k^{r-2}N'_k \backslash H_k^{r-1}$. Certainly $nu \rightarrow \psi'_1(u)\psi_1(n)$ gives a character ψ_2 on NN' , which is the unipotent radical $N^{r-2,1,1}$ of the $(r-2, 1, 1)$ parabolic. Thus, so far,

$$f(g) = \sum_{\gamma \in \Theta_k^{r-2}N'_k \backslash H_k} \int_{N_k^{r-2,1,1} \backslash N_\mathbb{A}^{r-2,1,1}} \bar{\psi}_2(n) f(n\gamma g) dn$$

We need a separate notation for unipotent radicals inside $H = H^{r-1} \approx GL_{r-1}$: let U^{b_1, \dots, b_m} be the unipotent radical of the standard parabolic of H with blocks of size b_1, \dots, b_m along the diagonal. Then

$$\Theta^{r-2} \cdot N' = H^{r-3} \cdot U^{r-3,1,1}$$

Thus,

$$f(g) = \sum_{\gamma \in H_k^{r-3}U_k^{r-3,1,1} \backslash H_k} \int_{N_k^{r-2,1,1} \backslash N_\mathbb{A}^{r-2,1,1}} \bar{\psi}_2(n) f(n\gamma g) dn$$

We repeat the induction step once more. For each $\gamma \in H_k^{r-3}U_k^{r-3,1,1} \backslash H_k$, the function

$$h \rightarrow \int_{N_k^{r-2,1,1} \backslash N_\mathbb{A}^{r-2,1,1}} \bar{\psi}_2(n) f(nh\gamma g) dn \quad (\text{for } h \in H^{r-2})$$

is left invariant under N'_k , where

$$N' = \{u_x = \begin{pmatrix} 1_{r-3} & x & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \text{ with } x = (r-3)\text{-by-}1\}$$

For each γ , there is a Fourier expansion along N' ,

$$\int_{N_k^{r-2,1,1} \backslash N_\mathbb{A}^{r-2,1,1}} \bar{\psi}_2(n) f(nh\gamma g) dn = \sum_{\psi'} \int_{N'_k \backslash N'_\mathbb{A}} \bar{\psi}'(u) \int_{N_k^{r-2,1,1} \backslash N_\mathbb{A}^{r-2,1,1}} \bar{\psi}_2(n) f(nuh\gamma g) dn du$$

The summand for trivial ψ' is 0, because f is a cuspform, and

$$N' \cdot (\ker \psi_2 \text{ on } N^{r-2,1,1}) \supset \left\{ \begin{pmatrix} 1_{r-3} & * & * & * \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \right\} = \text{unipotent radical of } (r-3, 3) \text{ parabolic}$$

The rational points of

$$H^{r-3} = \left\{ \begin{pmatrix} GL_{r-3} & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \right\}$$

act transitively on non-trivial characters ψ' of $U_k \backslash U_{\mathbb{A}}$. Let $\psi'_2(u_x) = \psi_o(x_{r-3})$. The isotropy group of ψ'_2 is

$$\Theta^{r-3} = \left\{ \begin{pmatrix} GL_{r-4} & * & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \right\}$$

Note that

$$\Theta^{r-3} \cdot U^{r-3,1,1} = H^{r-4} \cdot U^{r-4,1,1,1}$$

Setting $h = 1$,

$$f(g) = \sum_{\gamma \in H_k^{r-3} U_k^{r-3,1,1} \backslash H_k} \sum_{\delta \in \Theta_k^{r-3} \backslash H_k^{r-3}} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'_2(\delta u \delta^{-1}) \int_{N_k^{r-2,1,1} \backslash N_{\mathbb{A}}^{r-2,1,1}} \bar{\psi}_2(n) f(nu\gamma g) dn du$$

Replace u by $\delta^{-1}u\delta$, and n by $\delta^{-1}n\delta$, noting that conjugation by $N'_{\mathbb{A}}$ leaves ψ_2 invariant:

$$f(g) = \sum_{\gamma \in H_k^{r-3} U_k^{r-3,1,1} \backslash H_k} \sum_{\delta \in \Theta_k^{r-3} \backslash H_k^{r-3}} \int_{N'_k \backslash N'_{\mathbb{A}}} \bar{\psi}'_2(u) \int_{N_k^{r-2,1,1} \backslash N_{\mathbb{A}}^{r-2,1,1}} \bar{\psi}_2(n) f(nu\delta\gamma g) dn du$$

The double sum over $\delta\gamma$ can be regrouped into a single sum of $\gamma \in H_k^{r-4} \cdot U_k^{r-4,1,1,1} \backslash H_k$. Let

$$\psi_3 \begin{pmatrix} 1_{r-4} & 0 & * & * & * \\ 0 & 1 & x_{r-3,r-2} & * & * \\ 0 & 0 & 1 & x_{r-2,r-1} & * \\ 0 & 0 & 0 & 1 & x_{r-1,r} \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} = \psi_o(x_{r-3,r-2} + x_{r-2,r-1} + x_{r-1,r})$$

Then

$$f(g) = \sum_{\gamma \in H_k^{r-4} U_k^{r-4,1,1,1} \backslash H_k} \int_{N_k^{r-3,1,1,1} \backslash N_{\mathbb{A}}^{r-2,1,1,1}} \bar{\psi}_3(n) f(n\gamma g) dn$$

By induction, with U^{\min} the unipotent radical of the standard minimal parabolic in H , and N^{\min} the unipotent radical of the standard minimal parabolic in G ,

$$f(g) = \sum_{\gamma \in U_k^{\min} \backslash H_k} \int_{N_k^{\min} \backslash N_{\mathbb{A}}^{\min}} \bar{\psi}_{\text{std}}(n) f(n\gamma g) dn$$

Letting the Whittaker function attached to f be

$$W_f(g) = \int_{N_k^{\min} \backslash N_{\mathbb{A}}^{\min}} \bar{\psi}_{\text{std}}(n) f(n g) dn$$

the Fourier expansion is

$$f(g) = \sum_{\gamma \in U_k^{\min} \backslash H_k} W_f(\gamma g)$$

2. The Hecke-type integral representation: $GL_n \times GL_{n-1}$

Still let H denote the copy of GL_{r-1} in the standard Levi component of the standard $(r-1, 1)$ parabolic of $G = GL_r$. For cuspform f on GL_r and cuspform F on $H \approx GL_{r-1}$, the Hecke-type integral representation

$$\int_{H_k \backslash H_{\mathbb{A}}} |\det h|^{s-\frac{1}{2}} f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix} F(h) dh$$

produces the L -function $\Lambda(s, f \times F)$, up to normalization, as follows. Expressing f in its Fourier expansion, as above, unwind:

$$\int_{H_k \backslash H_{\mathbb{A}}} |\det h|^{s-\frac{1}{2}} \sum_{\gamma \in U_k^{\min} \backslash H_k} W_f(\gamma \cdot \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix}) F(h) dh = \int_{U_k^{\min} \backslash H_{\mathbb{A}}} |\det h|^{s-\frac{1}{2}} W_f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix} F(h) dh$$

The function $|\det h|^{s-\frac{1}{2}} W_f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix}$ is left ψ_{std} -equivariant under $U_{\mathbb{A}}^{\min}$, so rewrite

$$\begin{aligned} & \int_{U_k^{\min} \backslash H_{\mathbb{A}}} |\det h|^{s-\frac{1}{2}} W_f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix} F(h) dh \\ &= \int_{U_{\mathbb{A}}^{\min} \backslash H_{\mathbb{A}}} |\det h|^{s-\frac{1}{2}} W_f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix} \left(\int_{U_k^{\min} \backslash U_{\mathbb{A}}^{\min}} \psi_{\text{std}}(u) F(uh) du \right) dh \\ &= \int_{U_{\mathbb{A}}^{\min} \backslash H_{\mathbb{A}}} |\det h|^{s-\frac{1}{2}} W_f \begin{pmatrix} h & 0 \\ 0 & 1 \end{pmatrix} W_F(h) dh \end{aligned}$$

where we note that the Whittaker function for F is formed with the complex-conjugated character ψ_{std} restricted from N^{\min} to U^{\min} .

Under various hypotheses on f, F , the Whittaker functions factor over primes, and, then the zeta integral factors over primes, giving an *Euler product*

$$Z(s, f \times F) = \prod_v \left(\int_{U_v^{\min} \backslash H_v} |\det h_v|^{s-\frac{1}{2}} W_{f,v} \begin{pmatrix} h_v & 0 \\ 0 & 1 \end{pmatrix} W_{F,v}(h_v) dh_v \right)$$

Further, at places v where f and F are *spherical*, via an Iwasawa decomposition $H = U_v M_v^{\min} K_v$ with M^{\min} the standard Levi component of the minimal parabolic in H , the v^{th} local integral becomes a much smaller, $(r-1)$ -dimensional integral:

$$\int_{M_v^{\min}} |\det m_v|^{s-\frac{1}{2}} W_{f,v} \begin{pmatrix} m & 0 \\ 0 & 1 \end{pmatrix} W_{F,v}(m) \frac{dm}{\delta(m)} \quad \left(\text{with } m = \begin{pmatrix} m_1 & & 0 \\ & \ddots & \\ 0 & & m_{r-1} \end{pmatrix} \in GL_{r-1} \right)$$

where $\delta(m)$ is the modular function of M_v^{\min} on U_v^{\min} .

[2.0.1] **Remark:** The most important normalization constants are $\rho_f(1)$ and $\rho_F(1)$, which are the (higher-rank analogues of) leading Fourier coefficients of f and F , when f and F are normalized to have L^2 -norm 1. Further, it is less clear that the archimedean local zeta integral is the correct gamma factor. Thus, even with everywhere-spherical f and F ,

$$\text{zeta-integral } Z(s, f \times F) = \rho_f(1) \cdot \rho_{\overline{F}}(1) \cdot (\text{archimedean integrals}) \cdot L(s, f \times F)$$

[2.0.2] **Remark:** The analytic continuation of this Euler product follows from the original integral representation, with relatively straightforward estimates on the cuspforms f, F . The functional equation comes essentially from replacing h by h -transpose-inverse in the integral. The effect of transpose-inverse on the local representations, hence, on the Euler factors, requires some further attention.

3. The Rankin-Selberg case $GL_n \times GL_n$

Now we need an auxiliary (degenerate) Eisenstein series

$$E_s(g) = \sum_{\gamma \in P_k^{n-1,1} \backslash GL_n(k)} \varphi_s(\gamma \cdot g)$$

where $P^{n-1,1}$ has Levi component $M \approx GL_{n-1} \times GL_1$, and

$$\varphi_s = \bigotimes_v \varphi_{s,v}$$

with $\varphi_{s,v}$ in a (degenerate) induced representation from $P_v^{n-1,1}$. Specifically, the representation induced from M should be of the form

$$m = \begin{pmatrix} A & 0 \\ 0 & d \end{pmatrix} \longrightarrow \left| \frac{\det A}{d^{n-1}} \right|^s \cdot \chi \left(\frac{\det A}{d^{n-1}} \right) \quad (\text{for } A \in GL_{n-1} \text{ and } d \in GL_1)$$

where $s \in \mathbb{C}$ and χ is a Hecke character. [4] We suppress reference to other data specifying the vector φ_s in the induced representation, although in practice this data must be chosen to accommodate bad-prime aspects of f, F , for example. The character χ must be chosen so that the central character of $E_s \cdot f \cdot \overline{F}$ is trivial, or the following integral is not well-defined.

The Rankin-Selberg zeta integral attached to two cuspforms f, F on GL_n is [5]

$$Z(s, f \otimes \overline{F}) = \int_{Z_{\mathbb{A}} GL_n(k) \backslash GL_n(\mathbb{A})} E_s(g) \cdot f(g) \cdot \overline{F}(g) dg$$

Let H be the GL_{n-1} factor of the Levi component M . In the region of convergence, unwind the zeta integral by unwinding the Eisenstein series:

$$Z(s, f \otimes F) = \int_{Z_{\mathbb{A}} H_k N_k^{n-1,1} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot f(g) \cdot \overline{F}(g) dh$$

Let U^{\min} be the unipotent radical of the standard minimal parabolic in H . Expanding f in its Fourier-Whittaker expansion and unwinding, using the H_k -invariance of φ_s ,

$$\begin{aligned} Z(s, f \otimes F) &= \int_{Z_{\mathbb{A}} H_k N_k^{n-1,1} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot \sum_{\gamma \in U_k \backslash H_k} W_f(\gamma g) \cdot \overline{F}(g) dg \\ &= \int_{Z_{\mathbb{A}} U^{\min}_k N_k^{n-1,1} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot W_f(g) \cdot \overline{F}(g) dh = \int_{Z_{\mathbb{A}} N_k^{\min} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot W_f(g) \cdot \overline{F}(g) dh \end{aligned}$$

[4] The case of trivial χ is already useful. On the other hand, the character $|\cdot|^s$ can be incorporated into χ , if desired. Nevertheless, for analytical purposes, it is often convenient to separate the continuous parameter s from a parametrization of the compact part of the idele-class group: \mathbb{J}^1/k^\times where \mathbb{J}^1 is ideles of idele-norm 1.

[5] The complex conjugation on F avoids certain uninteresting technicalities, as will become apparent.

where $N^{\min} = U^{\min} N^{n-1,1}$ is the unipotent radical of the standard minimal parabolic P^{\min} in GL_n . Since φ_s is left $N_{\mathbb{A}}$ -invariant, and the Whittaker function W_f is left $N_{\mathbb{A}}$, ψ_{std} -equivariant, with the standard character ψ_{std} on $N_{\mathbb{A}}$,

$$\begin{aligned} \int_{Z_{\mathbb{A}} N_k^{\min} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot W_f(g) \cdot \overline{F}(g) \, dh &= \int_{Z_{\mathbb{A}} N_{\mathbb{A}}^{\min} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot W_f(g) \left(\int_{N_k^{\min} \backslash N_{\mathbb{A}}^{\min}} \psi_{\text{std}}(u) \cdot \overline{F}(ug) \, du \right) \, dh \\ &= \int_{Z_{\mathbb{A}} N_{\mathbb{A}}^{\min} \backslash GL_n(\mathbb{A})} \varphi_s(g) \cdot W_f(g) \cdot \overline{W_F}(g) \, dg \end{aligned}$$

It is here that the pre-emptive complex-conjugation of F gives the Whittaker function of F with respect to ψ_{std} , rather than with respect to its complex conjugate.

Under various hypotheses, the Whittaker functions factor over primes. When we take φ_s to be a monomial tensor, we have an Euler product

$$Z(s, f \otimes \overline{F}) = \prod_v \int_{Z_v N_v^{\min} \backslash GL_n(k_v)} \varphi_s(g) \cdot W_{f,v}(g) \cdot \overline{W_{F,v}}(g) \, dg \quad (\text{with inducing data suppressed})$$

[3.0.1] Remark: The most important normalization constants are $\rho_f(1)$ and $\rho_F(1)$, the (higher-rank analogues of) leading Fourier coefficients of f and F , when f and F are normalized to have L^2 -norm 1. It is less clear that the archimedean local zeta integral is the correct gamma factor. Thus, even with everywhere-spherical f and F ,

$$\text{zeta-integral } Z(s, f \times F) = \rho_f(1) \cdot \rho_{\overline{F}}(1) \cdot (\text{archimedean integrals}) \cdot L(s, f \times F)$$

[3.0.2] Remark: The analytic continuation of the zeta integral follows from the original integral representation, with relatively straightforward estimates on the cuspforms f, F , and from the meromorphic continuation and function equation of E_s . For this very degenerate Eisenstein series, the analytic continuation and functional equation follow from Poisson summation.

4. Comments on $GL_m \times GL_n$ with $m \leq n - 2$

As Jacquet, Piatetski-Shapiro, and Shalike found, for $m < n - 1$ an intermediate integration is necessary. In the literature, such auxiliary integrations are often called *unipotent* integrations, and occur in other situations, as well. In the end, one has a zeta integral

$$Z(s, f \otimes F) = \int_{GL_m(k) \backslash GL_m(\mathbb{A})} |\det h|^{s-\frac{1}{2}} \cdot (\text{proj}_m^n f) \left(\begin{pmatrix} h & 0 \\ 0 & 1_{n-m} \end{pmatrix} \right) \cdot F(h) \, dh$$

where the projection operator proj_m^n is the identity map for $m = n - 1$, but non-trivial otherwise, described as follows. Let

$$N = N_m^r = \begin{pmatrix} 1_{m+1} & * & \cdots & * \\ & 1 & & \vdots \\ & & \ddots & * \\ & & & 1 \end{pmatrix} = \text{unipotent radical of } (m+1, 1, 1, \dots, 1) \text{ parabolic}$$

The proper definition of the projection turns out to be

$$(\text{proj}_m^r f)(g) = |\det h|^{\frac{r-(m+1)}{2}} \int_{N_k \backslash N_{\mathbb{A}}} \overline{\psi}(n) f(n g) \, dn$$

For example, for $r = 3$ and $m = 1$, with F trivial on GL_1 , the standard L -function attached to f is essentially the zeta integral

$$\begin{aligned} Z(s, f) &= \int_{GL_1(k) \backslash GL_1(\mathbb{A})} |h|^{s - \frac{1}{2} + \frac{3 - (1+1)}{2}} \text{proj}_1^3 f \begin{pmatrix} h & & \\ & 1 & \\ & & 1 \end{pmatrix} dh \\ &= \int_{GL_1(k) \backslash GL_1(\mathbb{A})} |h|^s \int_{k \backslash \mathbb{A}} \int_{k \backslash \mathbb{A}} \bar{\psi}_o(y) f \left(\begin{pmatrix} 1 & x \\ & 1 & y \\ & & 1 \end{pmatrix} \begin{pmatrix} h & & \\ & 1 & \\ & & 1 \end{pmatrix} \right) dx dy dh \end{aligned}$$

See [Cogdell 2003,07,08] for many further details about this general case.

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