

Harmonic analysis, on \mathbb{R} , \mathbb{R}/\mathbb{Z} , \mathbb{Q}_p , \mathbb{A} , and \mathbb{A}_k/k , key ingredients in Iwasawa-Tate:

Characters and *Fourier transforms* on \mathbb{R} , \mathbb{Q}_p , \mathbb{A} , and for completions and adèles of number fields.

Fourier inversion expresses nice functions as *superpositions* (integrals) of characters.

Schwartz spaces $\mathcal{S}(\mathbb{R})$, $\mathcal{S}(\mathbb{Q}_p)$, $\mathcal{S}(\mathbb{A})$ of very-nice functions on \mathbb{R} , \mathbb{Q}_p , \mathbb{A} , and k_v , \mathbb{A}_k for number fields k .

Adelic Poisson summation

$$\sum_{x \in k} f(x) = \sum_{x \in k} \hat{f}(x) \quad (\text{for } f \in \mathcal{S} \mathbb{A}_k)$$

Recap:

Corollary: Given *non-trivial* $\psi \in k_v^\vee$, every other element of k_v^\vee is of the form $x \rightarrow \psi(\xi \cdot x)$ for some $\xi \in k_v$. Similarly, given *non-trivial* $\psi \in \mathbb{A}_k^\vee$, every other element of \mathbb{A}_k^\vee is of the form $x \rightarrow \psi(\xi \cdot x)$ for some $\xi \in \mathbb{A}_k$.

The **standard character** ψ_1 on \mathbb{Q}_p is as follows: given $x \in \mathbb{Q}_p$, there is $x' \in p^{-k}\mathbb{Z}$ for some $k \in \mathbb{Z}$, such that $x - x' \in \mathbb{Z}_p$, and

$$\psi_1(x) = e^{-2\pi i x'} \quad (\text{sign choice for later})$$

For $\xi \in \mathbb{Q}_p$, let

$$\psi_\xi(x) = \psi_1(\xi \cdot x) \quad (\text{for } x, \xi \in \mathbb{Q}_p)$$

For a finite extension k_v of \mathbb{Q}_p , the **standard character** is

$$\psi_\xi(x) = \psi_1(\text{tr}_{\mathbb{Q}_p}^{k_v}(\xi \cdot x)) \quad (\text{for } x, \xi \in k_v)$$

Proposition (later): the compact-open topology makes \widehat{G} an abelian (locally-compact, Hausdorff) topological group.

Theorem: Unitary dual of a *compact* abelian group is *discrete*, unitary dual of a *discrete* abelian group is *compact*. ///

Fourier transforms, Fourier inversion, Schwartz spaces of functions, adelic Poisson summation Unsurprisingly, the Fourier transform on k_v is

$$\mathcal{F} f(\xi) = \widehat{f}(\xi) = \int_{k_v} \overline{\psi_\xi(x)} f(x) dx$$

where, given the characters ψ_ξ (for example, the standard ones), the Haar measures are normalized so that *Fourier inversion* holds exactly:

$$f(x) = \int_{k_v} \psi_\xi(x) \widehat{f}(\xi) d\xi \quad (\text{for nice functions } f)$$

Why does Fourier inversion hold at all?

The usual space $\mathcal{S}(\mathbb{R})$ of *Schwartz functions* on \mathbb{R} consists of infinitely-differentiable functions all of whose derivatives are of *rapid decay*, decaying more rapidly at $\pm\infty$ than every $1/|x|^N$. Its topology is given by semi-norms

$$\nu_{k,N}(f) = \sup_{0 \leq i \leq k} \sup_{x \in \mathbb{R}} \left((1 + |x|)^N \cdot |f^{(i)}(x)| \right)$$

for $0 \leq k \in \mathbb{Z}$ and $0 \leq N \in \mathbb{Z}$.

There are countably-many associated (pseudo-) metrics $d_{k,N}(f, g) = \nu_{k,N}(f - g)$, so $\mathcal{S}(\mathbb{R})$ is naturally *metrizable*.

The usual two-or-three-epsilon arguments show that $\mathcal{S}(\mathbb{R})$ is *complete* metrizable.

Theorem: \mathcal{F} is a topological isomorphism $\mathcal{S}(\mathbb{R}) \rightarrow \mathcal{S}(\mathbb{R})$.

Before describing the Schwartz space $\mathcal{S}(\mathbb{Q}_p)$ and proving Fourier inversion, sample computations of Fourier transforms are useful.

We need a simply-described function on \mathbb{Q}_p which is its own Fourier transform, to play a role analogous to that of the Gaussian in the archimedean case.

Claim: With Fourier transform on \mathbb{Q}_p defined via the standard character $\psi_1(x) = e^{-2\pi i x'}$ (where $x' \in p^{-\infty}\mathbb{Z}_p$ and $x - x' \in \mathbb{Z}_p$), the characteristic function of \mathbb{Z}_p is its own Fourier transform.

Proof: Let f be the characteristic function of \mathbb{Z}_p . Then

$$\widehat{f}(\xi) = \int_{\mathbb{Q}_p} \overline{\psi_\xi(x)} f(x) dx = \int_{\mathbb{Z}_p} \overline{\psi_1(\xi \cdot x)} dx = \int_{\mathbb{Z}_p} \psi_1(-\xi \cdot x) dx$$

Recall a form of the *cancellation lemma*: (a tiny case of *Schur orthogonality*...)

Lemma: Let $\psi : K \rightarrow \mathbb{C}^\times$ be a continuous group homomorphism on a compact group K . Then

$$\int_K \psi(x) dx = \begin{cases} \text{meas}(K) & (\text{for } \psi = 1) \\ 0 & (\text{for } \psi \neq 1) \end{cases}$$

Proof of Lemma: Yes, of course, the measure is a Haar measure on K . Since K is *compact*, it is *unimodular*.

For ψ trivial, of course the integral is the total measure of K .

For ψ non-trivial, there is $y \in K$ such that $\psi(y) \neq 1$. Using the invariance of the measure, change variables by replacing x by xy :

$$\int_K \psi(x) dx = \int_K \psi(xy) d(xy) = \int_K \psi(x) \psi(y) dx = \psi(y) \int_K \psi(x) dx$$

Since $\psi(y) \neq 1$, the integral is 0. ///

Apply the lemma to the integrals computing the Fourier transform of the characteristic function f of \mathbb{Z}_p . Giving the compact group \mathbb{Z}_p measure 1,

$$\widehat{f}(\xi) = \int_{\mathbb{Z}_p} \psi_1(-\xi \cdot x) dx = \begin{cases} 1 & (\psi_1(-\xi x) = 1 \text{ for } x \in \mathbb{Z}_p) \\ 0 & (\text{otherwise}) \end{cases}$$

On one hand, for $\xi \in \mathbb{Z}_p$, certainly $\psi_1(\xi x) = 1$ for $x \in \mathbb{Z}_p$. On the other hand, for $\xi \notin \mathbb{Z}_p$, there is $x \in \mathbb{Z}_p$ such that, for example, $\xi \cdot x = 1/p$. Then

$$\psi_1(-\xi \cdot x) = \psi_1\left(\frac{-1}{p}\right) = e^{+2\pi i \cdot \frac{1}{p}} \neq 1$$

Thus, ψ_ξ is not trivial on \mathbb{Z}_p , so the integral is 0. Thus, the characteristic function of \mathbb{Z}_p is its own Fourier transform. ///

Claim: With standard Fourier transform on \mathbb{Q}_p , the Fourier transform of the characteristic function of $p^k\mathbb{Z}_p$ is p^{-k} times the characteristic function of $p^{-k}\mathbb{Z}_p$.

Proof: Let f be the characteristic function of $p^k\mathbb{Z}_p$, so

$$\begin{aligned}\widehat{f}(\xi) &= \int_{\mathbb{Q}_p} \overline{\psi}_\xi(x) f(x) dx = \int_{p^k\mathbb{Z}_p} \overline{\psi}_1(\xi \cdot x) dx \\ &= |p^k|_p \cdot \int_{\mathbb{Z}_p} \psi_1(-\xi \cdot x/p^k) dx = p^{-k} \cdot \int_{\mathbb{Z}_p} \psi_1(-\xi \cdot x/p^k) dx\end{aligned}$$

This reduces to the previous computation: by *cancellation*, for $\xi/p^k \notin \mathbb{Z}_p$ the character $x \rightarrow \psi_1(-\xi x/p^k)$ is non-trivial, so the integral is 0. Otherwise, the integral is 1. ///

Claim: With standard Fourier transform on \mathbb{Q}_p , the Fourier transform of the characteristic function of $\mathbb{Z}_p + y$ is ψ_y times the characteristic function of \mathbb{Z}_p .

Proof: Let f be the characteristic function of $\mathbb{Z}_p + y$, so

$$\begin{aligned}\widehat{f}(\xi) &= \int_{\mathbb{Q}_p} \overline{\psi}_\xi(x) f(x) dx = \int_{\mathbb{Z}_p+y} \overline{\psi}_1(\xi \cdot x) dx \\ &= \int_{\mathbb{Z}_p} \psi_1(-\xi \cdot (x + y)) dx = \psi_1(-\xi \cdot y) \int_{\mathbb{Z}_p} \psi_1(-\xi \cdot x) dx \\ &= \psi_1(-\xi \cdot y) \cdot f(\xi)\end{aligned}$$

by the previous computation.

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Combining the two computations above,

$$\mathcal{F}\left(\text{char fcn } p^k\mathbb{Z}_p + y\right) = \psi_y \cdot p^{-k} \cdot (\text{char fcn } p^{-k}\mathbb{Z}_p)$$

Conveniently, products $\psi_y \cdot (\text{char fcn } p^{-k}\mathbb{Z}_p)$ are in the same class of functions, since ψ_y has a kernel which is an open (and compact) neighborhood of 0, so we *this class of functions is mapped to itself under Fourier transform*.

Recall the earlier lemma proving that these *special* simple functions consisting of finite linear combinations of characteristic functions of sets $p^k\mathbb{Z}_p + y$ are *dense* in $C_c^o(\mathbb{Q}_p)$.

The space of **Schwartz functions** $\mathcal{S}(\mathbb{Q}_p)$ on \mathbb{Q}_p is the vector space of these special simple functions, that is, finite linear combinations of characteristic functions of sets $p^k\mathbb{Z}_p + y$.

Yes, unlike the archimedean case, p -adic Schwartz functions are *compactly supported*.

p -adic Fourier inversion is the claim that

$$f(x) = \int_{\mathbb{Q}_p} \psi_\xi(x) \widehat{f}(\xi) d\xi \quad (\text{for } f \in \mathcal{S}(\mathbb{Q}_p))$$

Proof: We have essentially proven this in the computations above, if we keep track, as follows. Let f° be the characteristic function of \mathbb{Z}_p . We computed $\widehat{f^\circ} = f$. Let δ_t be the dilation operator $\delta_t f(x) = f(t \cdot x)$ for $t \in \mathbb{Q}_p^\times$. We computed, by changing variables in the integral defining the Fourier transform, that

$$\mathcal{F}(\delta_t f) = \frac{1}{|t|_p} \cdot \delta_{1/t}(\mathcal{F} f)$$

Let τ_y be the translation operator $\tau_y f(x) = f(x + y)$. By changing variables,

$$\mathcal{F}(\tau_y f) = \psi_y \cdot (\mathcal{F} f)$$

It is convenient to also compute that

$$\begin{aligned}\mathcal{F}(\psi_y \cdot f)(\xi) &= \int_{\mathbb{Q}_p} \bar{\psi}_\xi(x) \cdot \psi_y(x) f(x) dx \\ &= \int_{\mathbb{Q}_p} \bar{\psi}_{\xi-y}(x) f(x) dx = \widehat{f}(\xi - y) = \tau_{-y}(\mathcal{F} f)\end{aligned}$$

Let \mathcal{F}^* be the integral for Fourier inversion, namely,

$$\mathcal{F}^* f(x) = \int_{\mathbb{Q}_p} \psi_\xi(x) f(\xi) d\xi$$

Similar computations give

$$\mathcal{F}^*(\delta_t f) = \frac{1}{|t|_p} \delta_{1/t}(\mathcal{F}^* f) \quad \mathcal{F}^*(\tau_y f) = \psi_{-y}(\mathcal{F}^* f)$$

and

$$\mathcal{F}^*(\psi_y f) = \tau_y(\mathcal{F}^* f)$$

Since every element of $\mathcal{S}(\mathbb{Q}_p)$ is a linear combination of images of f° under dilation and translation, it suffices to give a sort of inductive proof of Fourier inversion:

$$\mathcal{F}^* \mathcal{F}(\tau_y f) = \mathcal{F}^* \psi_y \mathcal{F} f = \tau_y \mathcal{F}^* \mathcal{F} f$$

$$\mathcal{F}^* \mathcal{F}(\delta_t f) = \mathcal{F}^* \frac{1}{|t|_p} \delta_{1/t} \mathcal{F} f = \frac{1}{|t|_p} \frac{1}{|1/t|_p} \delta_t \mathcal{F}^* \mathcal{F} f = \delta_t \mathcal{F}^* \mathcal{F} f$$

Similarly for multiplication by ψ_y . Since $\mathcal{F}^* \mathcal{F} f^\circ = \mathcal{F}^* f^\circ = f^\circ$, we have Fourier inversion on $\mathcal{S}(\mathbb{Q}_p)$. ///

Remark: p -adic Fourier inversion is much easier than on \mathbb{R} .

The space $\mathcal{S}(\mathbb{A})$ of Schwartz functions on the adeles is finite linear combinations of *monomial* functions

$$\left(\bigotimes_{v \leq \infty} f_v \right) (\{x_v\}) = \prod_v f_v(x_v)$$

with $f_v \in \mathcal{S}(\mathbb{Q}_v)$, and where *for all but finitely-many* v the local function f_v is the characteristic function of \mathbb{Z}_v .

Fourier transform on $\mathcal{S}(\mathbb{A})$ is the product of all the local Fourier transforms, and Fourier inversion follows for $\mathcal{S}(\mathbb{A})$ because it holds for each $\mathcal{S}(\mathbb{Q}_v)$.

Remark: We do not directly need it, but one might reflect on what the natural topology is on $\mathcal{S}(\mathbb{Q}_p)$, especially to have it be *complete*.

Fourier series on \mathbb{A}/k : For a unimodular topological group G , let $L^2(G)$ be the *completion* of $C_c^o(G)$ with respect to the usual L^2 -norm given by

$$|f|^2 = \int_G |f(g)|^2 dg \quad (\text{for } f \in C_c^o(G))$$

Remark: The measurable-function $L^2(G)$ contains this completion, and is provably equal, but we only need integrals of continuous compactly-supported functions.

Theorem: For a compact abelian group G , with total measure 1, the continuous group homomorphisms (*characters*) $\psi : G \rightarrow \mathbb{C}^\times$ form an orthonormal *Hilbert-space basis* for $L^2(G)$. That is,

$$L^2(G) = \text{completion of } \bigoplus_{\psi \in G^\vee} \mathbb{C} \cdot \psi$$

Remark: A *Hilbert-space basis* of a Hilbert space V is not a *vector-space basis* for V , but only for a dense subspace.

