

**Iwasawa-Tate** on  $\zeta$ -functions and  $L$ -functions

After the main part, namely, *analytic continuation* and *functional equation* of global zeta integrals...

- [recap] Local functional equations

$$\frac{Z_v(s, f_v)}{Z_v(1-s, \widehat{f}_v)} = \frac{Z_v(s, g_v)}{Z_v(1-s, \widehat{g}_v)} \quad (\text{for } f, g \in \mathcal{S}(k_v))$$

- [recap] Unramified and ramified archimedean zeta integrals
- Convergence of local zeta integrals
- Archimedean Fourier transform eigenfunctions (Hecke's identity)

**Local functional equations:** For  $0 < \operatorname{Re}(s) < 1$ , the local zeta integrals are absolutely convergent, and

$$\frac{Z_v(s, \chi, f_v)}{Z_v(1-s, \bar{\chi}, \widehat{f}_v)} = \frac{Z_v(s, \chi, g_v)}{Z_v(1-s, \bar{\chi}, \widehat{g}_v)} \quad (\text{for } f, g \in \mathcal{S}(k_v))$$

The first point is that local zeta integrals with any Schwartz functions whatsoever are meromorphic, granting that *one* (non-zero) local zeta integral at  $v$  is meromorphic:

$$Z_v(1-s, \bar{\chi}, \widehat{g}) = \frac{Z_v(1-s, \bar{\chi}, \bar{f}) \cdot Z_v(s, \chi, g)}{Z_v(s, \chi, f)}$$

gives the meromorphic continuation of  $Z_v(s, \widehat{g})$ , etc.

**Remark:** The *local* functional equations do not yield the *global* functional equation! They only prove the essential irrelevance of the Schwartz data.

*Proof:* [Recap with  $\chi\dots$ ] Take  $0 < \operatorname{Re}(s) < 1$ , so local zeta integrals for both  $Z(s, \chi, f)$  and  $Z(1 - s, \bar{\chi}, \hat{g})$  converge absolutely. Direct computation: *expand the definition and change variables...*

Suppress all subscripts.

$$\begin{aligned} & Z(s, \chi, f) Z(1 - s, \bar{\chi}, \hat{g}) \\ &= \int_{k^\times} \int_{k^\times} |x|^s \chi(x) f(x) |y|^{1-s} \bar{\chi}(y) \hat{g}(y) d^\times x d^\times y \\ &= \int_{k^\times} \int_{k^\times} \int_k \bar{\psi}(y\eta) |x|^s \chi(x) f(x) |y|^{1-s} \bar{\chi}(y) g(\eta) d^+ \eta d^\times y d^\times x \end{aligned}$$

Replacing  $y$  by  $\frac{yx}{\eta}$ , this is

$$\begin{aligned} & \int_k \int_{k^\times} \int_{k^\times} \bar{\psi}(yx) |x|^s f(x) \left| \frac{yx}{\eta} \right|^{1-s} \bar{\chi}(y) \chi(\eta) g(\eta) d^\times y d^\times x d^+ \eta \\ &= \int_k \int_{k^\times} \int_{k^\times} \bar{\psi}(yx) |x|^1 f(x) \bar{\chi}(y) |y|^{1-s} \chi(\eta) |\eta|^{s-1} g(\eta) d^\times y d^\times x d^+ \eta \end{aligned}$$

The measure  $|x| \cdot d^\times x$  is a constant multiple of the *additive* Haar measure  $d^+ x$ . The constant is irrelevant, since it cancels itself in the rearrangement:

$$|x| |\eta|^{-1} d^\times x d^+ \eta = d^+ x d^\times \eta$$

Continuing,

$$\begin{aligned} & Z(s, \chi, f) Z(1 - s, \bar{\chi}, \hat{g}) \\ &= \int_{k^\times} \int_k \int_{k^\times} \bar{\psi}(yx) f(x) \bar{\chi}(y) |y|^{1-s} \chi(\eta) |\eta|^s g(\eta) d^\times y d^+ x d^\times \eta \\ &= \int_{k^\times} \int_{k^\times} \hat{f}(y) \bar{\chi}(y) |y|^{1-s} \chi(\eta) |\eta|^s g(\eta) d^\times y d^\times \eta \\ &= Z(1 - s, \bar{\chi}, \hat{f}) Z(s, \chi, g) \end{aligned}$$

This proves the local functional equation in  $0 < \operatorname{Re}(s) < 1$ . The general assertion follows from the Identity Principle. ///

**Complex zeta integrals [Recap]**

The product formula requires

$$|z|_{\mathbb{C}} = |N_{\mathbb{R}}^{\mathbb{C}}(z)|_{\mathbb{R}} = |z|^2 \quad (\text{the latter the } \textit{usual} \text{ norm})$$

Similarly, the *additive character*  $\psi_{\mathbb{C}}(z)$  is

$$\psi_{\mathbb{C}}(z) = \psi_{\mathbb{R}}(\text{tr}_{\mathbb{R}}^{\mathbb{C}}(z)) = e^{2\pi i(z+\bar{z})} = e^{4\pi i \text{Re}(z)}$$

The proper normalization of measure on  $\mathbb{C}$  for Iwasawa-Tate is

$$d_{\mathbb{C}}(z) = 2 \cdot d_{\text{usual}}(z)$$

The proper normalization of Gaussian to be its own Fourier transform is

$$f_{\mathbb{C}}(w) = e^{-2\pi w\bar{w}}$$

The standard *unramified* complex zeta integral is

$$\int_{\mathbb{C}} |z|_{\mathbb{C}}^s e^{-2\pi z\bar{z}} \frac{2 dz}{|z|_{\mathbb{C}}} = 2\pi \cdot (2\pi)^{-s} \Gamma(s)$$

The *ramified* unitary characters of  $\mathbb{C}^\times$  are

$$\chi_\ell(re^{i\theta})|re^{i\theta}|_{\mathbb{C}}^s = e^{i\ell\theta} \cdot r^{2s} \quad (\text{for } \ell \in \mathbb{Z})$$

The standard choice of Schwartz function for the complex zeta integral depends on the *sign* of  $\ell \in \mathbb{Z}$ . For  $\ell \geq 0$ , it is

$$\int_{\mathbb{C}} |z|_{\mathbb{C}}^s \chi_\ell(z) \bar{z}^\ell e^{-2\pi z\bar{z}} \frac{2 dz}{|z|_{\mathbb{C}}} = 2\pi \cdot (2\pi)^{-(s+\frac{\ell}{2})} \Gamma\left(s + \frac{\ell}{2}\right)$$

For  $\ell \leq 0$ , the standard ramified complex local zeta integral is

$$\begin{aligned}
& \int_{\mathbb{C}} |z|_{\mathbb{C}}^s \chi_{\ell}(z) |z|^{\ell} e^{-2\pi z \bar{z}} \frac{2 dz}{|z|_{\mathbb{C}}} \\
&= \int_0^{2\pi} \int_0^{\infty} r^{2s} e^{i\ell\theta} (re^{i\theta})^{-\ell} e^{-2\pi r^2} \frac{2 dz}{|z|_{\mathbb{C}}} \\
&= 4\pi \int_0^{\infty} r^{2s-\ell} e^{-2\pi r^2} r \frac{dr}{r^2} = 4\pi \int_0^{\infty} r^{2s-\ell} e^{-2\pi r^2} \frac{dr}{r} \\
&= 2\pi \int_0^{\infty} r^{s-\frac{\ell}{2}} e^{-2\pi r} \frac{dr}{r} = 2\pi \cdot (2\pi)^{-(s-\frac{\ell}{2})} \int_0^{\infty} r^{s-\frac{\ell}{2}} e^{-r} \frac{dr}{r} \\
&= 2\pi \cdot (2\pi)^{-(s-\frac{\ell}{2})} \Gamma\left(s - \frac{\ell}{2}\right) \quad (\text{for } \ell \leq 0)
\end{aligned}$$

Thus, for *both*  $\ell \geq 0$  and  $\ell \leq 0$ , the local zeta integral is

$$2\pi \cdot (2\pi)^{-(s+\frac{|\ell|}{2})} \Gamma\left(s + \frac{|\ell|}{2}\right) \quad (\text{for both } \ell \geq 0 \text{ and } \ell \leq 0)$$

**Convergence of local zeta integrals in  $\operatorname{Re} s > 0$ :**

As usual, suppose  $\chi_v$  is *unitary* meaning  $|\chi_v| = 1$ , since any non-unitary part could be absorbed into  $|\cdot|^s$ . Use the standard notation  $\sigma = \operatorname{Re}(s)$ .

Treat the non-archimedean case first. Since  $f_v \in \mathcal{S}(k_v)$ , for some  $n$  the support of  $f$  is contained in  $\varpi^{-n}\mathfrak{o}_v$ . Since  $f$  is locally constant and compactly supported, it has a finite bound  $C$ . Then

$$\begin{aligned} |Z_v(s, \chi_v, f_v)| &\leq \int_{k_v^\times} |x|^\sigma |\chi(x)| |f(x)| dx \\ &\leq C \cdot \int_{k_v^\times \cap \varpi^{-n}\mathfrak{o}_v} |x|^\sigma dx = C \cdot \int_{(k_v^\times \cap \varpi^{-n}\mathfrak{o}_v)/\mathfrak{o}_v^\times} |x|^\sigma \left( \int_{\mathfrak{o}_v^\times} 1 \right) dx \\ &= C \cdot \sum_{\ell=-n}^{\infty} |\varpi^\ell|_v^\sigma = C \cdot \frac{q_v^{n\sigma}}{1 - q_v^{-\sigma}} < \infty \quad (\text{for } \sigma = \operatorname{Re}(s) > 0) \end{aligned}$$



For  $k_v \approx \mathbb{R}$ , given  $f \in \mathcal{S}(\mathbb{R})$  for each  $N$  that

$$|f(x)| \ll_N (1 + |x|^2)^{-N}$$

With  $\sigma = \operatorname{Re}(s) > 0$ , the local zeta integral is

$$\begin{aligned} |Z_v(s, \chi_v, f_v)| &\ll_N \int_{\mathbb{R}^\times} |x|^\sigma |\chi_v(x)| (1 + x^2)^{-N} \frac{dx}{|x|} \\ &\ll \int_0^\infty |x|^{\sigma-1} (1 + x^2)^{-N} dx \\ &\ll \int_0^1 |x|^{\sigma-1} dx + \int_1^\infty |x|^{\sigma-1-2N} dx \end{aligned}$$

Given  $\sigma > 0$ , take  $N$  large enough so that  $\sigma - 1 - 2N < -1$  gives convergence. ///

Convergence of the complex integrals is similar...

## Archimedean Fourier transform eigenfunctions (Hecke's identity)

Whenever possible, we want local Schwartz functions  $f_v$  which are eigenfunctions for Fourier transform, preferably unchanged.

We want this so that in the *global* functional equation  $Z(s, \chi, f) = Z(1 - s, \bar{\chi}, \hat{f})$  as many local factors as possible are the same on both sides, apart from  $s \leftrightarrow 1 - s$  and  $\chi \leftrightarrow \bar{\chi}$ .

For absolutely unramified  $k_v/\mathbb{Q}_p$ , the characteristic function of the local integers  $\mathfrak{o}_v$  is its own Fourier transform. For  $\chi_v$  unramified at  $v$ , this gives the desired symmetry.

On  $\mathbb{R}$ , the Gaussian  $e^{-\pi x^2}$  is its own Fourier transform.

On  $\mathbb{R}$ , for ramified  $\chi_{\mathbb{R}}$ , that is, for  $\text{sgn}(x)|x|^s$ , the function  $f_v(x) = xe^{-\pi x^2}$  is multiplied by  $-i$  under Fourier transform.

*Proof:* One argument is by contour-shifting:

$$\begin{aligned}
 \widehat{f}_v(\xi) &= \int_{\mathbb{R}} e^{-2\pi i \xi x} x e^{-\pi x^2} dx \\
 &= \int_{\mathbb{R}} e^{-2\pi i \xi(x-i\xi)} (x - i\xi) e^{-\pi(x-i\xi)^2} dx \\
 &= e^{-\pi \xi^2} \int_{\mathbb{R}} (x - i\xi) e^{-\pi x^2} dx = e^{-\pi \xi^2} \cdot (0 - i\xi) = -i\xi e^{-\pi \xi^2}
 \end{aligned}$$

Done.

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**Remark:** While contour-shifting arguments have certain virtues, their computation-intensive, somewhat non-conceptual nature is not as helpful as we might hope.

Fourier transforms of Schwartz functions for ramified characters on  $k_v \approx \mathbb{C}$  are the critical sub-case of *Hecke's identity* on  $\mathbb{R}^n$ .

With, *usual* normalizations:

**Claim:** The Schwartz function  $(x \pm iy)^\ell e^{-\pi(x^2+y^2)}$  is an eigenfunction for Fourier transform, with eigenvalue  $i^{-\ell}$ .

*Proof:* Just do the case  $(x + iy)^\ell e^{-\pi(x^2+y^2)}$ . Rewrite this as  $z^\ell e^{-\pi z \bar{z}}$ , and rewrite the Fourier transform as

$$\begin{aligned} & \int_{\mathbb{C}} e^{-\pi i(z\bar{w} + \bar{z}w)} z^\ell e^{-\pi z \bar{z}} dz \\ &= (-\pi i)^{-\ell} \left( \frac{\partial}{\partial \bar{w}} \right)^\ell \int_{\mathbb{C}} e^{-\pi i(z\bar{w} + \bar{z}w)} e^{-\pi z \bar{z}} dz = (-\pi i)^{-\ell} \left( \frac{\partial}{\partial \bar{w}} \right)^\ell e^{-\pi w \bar{w}} \\ &= (-\pi i)^{-\ell} (-\pi w)^\ell e^{-\pi w \bar{w}} = i^{-\ell} \cdot w^\ell e^{-\pi w \bar{w}} \end{aligned}$$

This presumes  $\partial/\partial \bar{w}$  works as expected, which it does. ///

**Hecke's identity:** Let  $P$  be a homogeneous, degree  $d$  harmonic polynomial on  $\mathbb{R}^n$ , meaning that  $\Delta P = 0$ , where  $\Delta = \sum_j \partial^2 / \partial x_j^2$  is the usual Laplacian. Let  $\langle x, \xi \rangle = \sum_j x_j \xi_j$  be the usual pairing. Then  $P(x) e^{-\pi|x|^2}$  is a Fourier transform eigenfunction with eigenvalue  $i^{-d}$ :

$$\left( P(x) e^{-\pi|x|^2} \right)^\wedge(\xi) = i^{-d} \cdot P(\xi) e^{-\pi|\xi|^2}$$

[Proof later.]

**Remark:** Any other non-degenerate pairing  $x \times \xi \rightarrow \langle x, \xi \rangle$  and suitable associated operator  $\Delta$  works the same way.

**Remark:** Since  $f^\wedge(x) = f(-x)$ , necessarily  $f^{\wedge\wedge\wedge} = f$ , for all  $f$ . Thus, the only possible eigenvalues of Fourier transform are  $\pm 1, \pm i$ . Further, the corresponding components are easy to pick out: with  $\varepsilon \in \{\pm 1, \pm i\}$ , the  $\varepsilon^{th}$  component of  $f$  is

$$f + \varepsilon^{-1} \hat{f} + \varepsilon^{-2} \hat{\hat{f}} + \varepsilon^{-3} \hat{\hat{\hat{f}}}$$