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## Examples discussion 04

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[This document is  
http://www.math.umn.edu/~garrett/m/real/examples\_2017-18/real-disc-04.pdf]

[04.1] Fix  $x_o \in [a, b]$ . Show that  $\lambda(f) = f(x_o)$  is a continuous linear functional on  $C^o[a, b]$ .

**Discussion:** Recall that, for linear maps, continuity is equivalent to continuity at 0, which is equivalent to being *bounded*, in the sense that there exists a constant  $C$  such that  $|\lambda(f)| \leq C \cdot \|f\|_{C^o}$  for all  $f$ . Here,

$$|\lambda(f)| = |f(x_o)| \leq \sup_{x \in [a, b]} |f(x)| = \|f\|_{C^o}$$

so the constant  $C = 1$  succeeds. ///

[04.2] Prove that Cesaro summation

$$b_1 = \frac{a_1}{1}, \quad b_2 = \frac{a_1 + a_2}{2}, \quad b_3 = \frac{a_1 + a_2 + a_3}{3}, \dots$$

converts every convergent sequence  $a_1, a_2, \dots$  to a convergent sequence  $b_1, b_2, \dots$  with the same limit.

**Discussion:** Let  $\{a_n\}$  converge to  $A$ . Thus, given  $\varepsilon > 0$ , there is  $n_o$  be such for  $n > n_o$  we have  $|a_n - A| < \varepsilon$ . Let  $M = \max_{n \leq n_o} |a_n|$ . For  $n \geq n_o$ , by the triangle inequality,

$$\begin{aligned} |b_n - A| &= \frac{|(a_1 - A) + \dots + (a_n - A)|}{n} \leq \frac{|a_1| + \dots + |a_{n_o}|}{n} + \frac{n_o \cdot |A|}{n} + \frac{|a_{n_o+1} - A| + \dots + |a_n - A|}{n} \\ &< \frac{n_o \cdot M}{n} + \frac{n_o \cdot |A|}{n} + \varepsilon \end{aligned}$$

For  $n$  sufficiently large, depending on  $A, n_o$  and  $M$ , the sum of the first two terms can be made smaller than  $\varepsilon$ . Replace  $\varepsilon$  by  $\varepsilon/2$  throughout, if desired. ///

[04.3] (Collecting Fourier transform pairs...) Compute the Fourier transforms of

$$\chi_{[a, b]} \quad e^{-\pi x^2} \quad f(x) = \begin{cases} e^{-x} & (\text{for } x > 0) \\ 0 & (\text{for } x \leq 0) \end{cases}$$

**Discussion:** The first of these is direct:

$$\widehat{\chi_{[a, b]}}(\xi) = \int_{\mathbb{R}} e^{-2\pi i \xi x} \chi_{[a, b]}(x) dx = \int_a^b e^{-2\pi i \xi x} dx = \begin{cases} \frac{e^{-2\pi i \xi b} - e^{-2\pi i \xi a}}{-2\pi i \xi} & (\text{for } \xi \neq 0) \\ b - a & (\text{for } \xi = 0) \end{cases}$$

Since the latter function is *not* in  $L^1(\mathbb{R})$ , but *is* in  $L^2(\mathbb{R})$ , we define its Fourier transform (or inverse Fourier transform) *indirectly*, via either the inversion theorem, or by extending-by-continuity via Plancherel, expressing the function as an  $L^2$  limite of  $L^1$  functions.

The third is similarly direct:

$$\widehat{f}(\xi) = \int_0^\infty e^{-2\pi i \xi x} e^{-x} dx = \int_0^\infty e^{-(2\pi i \xi + 1)x} dx = \left[ \frac{e^{-(2\pi i \xi + 1)x}}{-(2\pi i \xi + 1)} \right]_0^\infty = \frac{1}{2\pi i \xi + 1}$$

Again, the latter function is not in  $L^1$ , but is in  $L^2$ , so its Fourier transform is most conveniently defined indirectly.

The Gaussian's Fourier transform is less trivial to evaluate, but is a very important example to have in hand, with many different applications throughout mathematics. One approach is as follows. Letting  $f(x) = e^{-\pi x^2}$ ,

$$\widehat{f}(\xi) = \int_{\mathbb{R}} e^{-2\pi i \xi x} e^{-\pi x^2} dx = \int_{\mathbb{R}} e^{-\pi(x^2 + 2i\xi x)} dx = \int_{\mathbb{R}} e^{-\pi(x^2 + i\xi)^2 - \pi\xi^2} dx = e^{-\pi\xi^2} \int_{\mathbb{R}} e^{-\pi(x+i\xi)^2} dx$$

by completing the square. The unobvious claim is that the integral does not depend on  $\xi$ , and, in fact, has value 1. Perhaps the optimal approach here is to observe that the integral is equal to a complex contour integral:

$$\int_{\mathbb{R}} e^{-\pi(x^2 + i\xi)^2} dx = \int_{i\xi - \infty}^{i\xi + \infty} e^{-\pi z^2} dz$$

along the line  $\text{Im}(z) = i\xi$ . Given the good decay of the integrand as  $|\text{Re}(z)| \rightarrow \infty$ , by Cauchy-Goursat theory, the contour can be *moved* to integration along the real line, giving

$$\int_{\mathbb{R}} e^{-\pi(x^2 + i\xi)^2} dx = \int_{i\xi - \infty}^{i\xi + \infty} e^{-\pi z^2} dz = \int_{-\infty}^{\infty} e^{-\pi x^2} dx = 1$$

The fact that the latter integral has value 1 comes from the usual trick involving polar coordinates:

$$\left( \int_{-\infty}^{\infty} e^{-\pi x^2} dx \right)^2 = \int_{\mathbb{R}^2} e^{-\pi(x^2 + y^2)} dx dy = \int_0^{2\pi} \int_0^{\infty} e^{-\pi r^2} r dr d\theta = 2\pi \int_0^{\infty} e^{-\pi r^2} r dr$$

Replacing  $r$  by  $\sqrt{t}$ , this is

$$\pi \int_0^{\infty} e^{-\pi t} dt = \pi \cdot \frac{1}{\pi} = 1$$

Thus, with the present normalization of Fourier transform and corresponding normalization of Gaussian, the Gaussian is its own Fourier transform. ///

[04.4] Show that  $\chi_{[a,b]} * \chi_{[c,d]}$  is a piecewise-linear function, and express it explicitly.

**Discussion:** Once enunciated, this fact (and the explicit expression) should be just a matter of book-keeping. We do assume that  $a \leq b$  and  $c \leq d$ . Also, by symmetry, without loss of generality we can suppose that  $|b - a| \geq |d - c|$ . This is used in the treatment of cases below.

$$\begin{aligned} (\chi_{[a,b]} * \chi_{[c,d]})(x) &= \int_{\mathbb{R}} \chi_{[a,b]}(x - y) \cdot \chi_{[c,d]}(y) dy = \int_c^d \chi_{[a,b]}(x - y) dy \\ &= \int_c^d \chi_{[a-x, b-x]}(-y) dy = \int_{-d}^{-c} \chi_{[x-b, x-a]}(y) dy = \text{meas}([-d, -c] \cap [x-b, x-a]) \end{aligned}$$

Looking at the cases of overlap, using  $b - a \geq d - c$ , this is

$$\left\{ \begin{array}{ll} 0 & \text{(for } x - a \leq -d, \text{ that is, } [x - b, x - a] \text{ is to the left of } [-d, -c]) \\ (x - a) - (-d) & \text{(for } x - b \leq -d \leq x - a \leq -c) \\ (-c) - (-d) & \text{(for } x - b \leq -d \leq -c \leq x - a, \text{ that is, } [-d, -c] \subset [x - b, x - a]) \\ (-c) - (x - b) & \text{(for } -d \leq x - b \leq -c \leq x - a) \\ 0 & \text{(for } x - b \geq -c, \text{ that is, } [x - b, x - a] \text{ is to the right of } [-d, -c]) \end{array} \right.$$

$$= \left\{ \begin{array}{ll} 0 & \text{(for } x \leq a - d) \\ x - a + d & \text{(for } a - d \leq x \leq a - c) \\ d - c & \text{(for } a - c \leq x \leq b - d) \\ b - c - x & \text{(for } b - d \leq x \leq b - c) \\ 0 & \text{(for } x \geq b - c) \end{array} \right.$$

We used the fact that  $b - a \geq d - c$  implies  $a - c \leq b - d$ . It is useful to consider the special configuration  $[a, b] = [-A, A]$  and  $[c, d] = [-B, B]$  with  $A \geq B \geq 0$ : the convolution is

$$\left\{ \begin{array}{ll} 0 & \text{(for } x \leq -A - B) \\ x + A + B & \text{(for } -A - B \leq x \leq -A + B) \\ 2B & \text{(for } -A + B \leq x \leq A - B) \\ A + B - x & \text{(for } A - B \leq x \leq A + B) \\ 0 & \text{(for } x \geq A + B) \end{array} \right.$$

In particular, the convolution is supported inside  $[-A - B, A + B]$ . Similarly, for  $f$  and  $g$  supported in  $[-a, a]$  and  $[-b, b]$ , the convolution is supported in  $[-a - b, a + b]$ . ///

[04.5] Evaluate the *Borwein integral*

$$\int_{\mathbb{R}} \frac{\sin x}{x} \cdot \frac{\sin x/3}{x/3} \cdot \frac{\sin x/5}{x/5} dx$$

**Discussion:** View this as an inner product and invoke Plancherel:

$$\int_{\mathbb{R}} \frac{\sin x}{x} \cdot \frac{\sin x/3}{x/3} \cdot \frac{\sin x/5}{x/5} dx = \left\langle \frac{\sin x}{x}, \frac{\sin x/3}{x/3} \cdot \frac{\sin x/5}{x/5} \right\rangle = \left\langle \left( \frac{\sin x}{x} \right)^{\widehat{\phantom{x}}}, \left( \frac{\sin x/3}{x/3} \cdot \frac{\sin x/5}{x/5} \right)^{\widehat{\phantom{x}}} \right\rangle$$

Since Fourier transform converts pointwise multiplication to convolution, this is

$$\left\langle \left( \frac{\sin x}{x} \right)^{\widehat{\phantom{x}}}, \left( \frac{\sin x/3}{x/3} \right)^{\widehat{\phantom{x}}} * \left( \frac{\sin x/5}{x/5} \right)^{\widehat{\phantom{x}}} \right\rangle$$

We have computed that

$$\chi_{[-a, a]}^{\widehat{\phantom{x}}}(\xi) = \frac{\sin 2\pi a \xi}{\pi \xi} = 2a \cdot \frac{\sin 2\pi a \xi}{2\pi a \xi}$$

That is, by linearity of Fourier transform,

$$\left(\frac{1}{2a}\chi_{[-a,a]}\right)^\wedge(\xi) = \frac{\sin(2\pi a)\xi}{(2\pi a)\xi}$$

By Fourier inversion, noting that  $\frac{\sin x}{x}$  is not in  $L^1$ , only in  $L^2$ , so the inverse transform is not necessarily the literal integral,

$$\left(\frac{\sin(2\pi a)\xi}{(2\pi a)\xi}\right)^\wedge(x) = \frac{1}{2a}\chi_{[-a,a]}(x)$$

Replacing  $a$  by  $a/2\pi$  gives

$$\left(\frac{\sin a\xi}{a\xi}\right)^\wedge(x) = \frac{\pi}{a}\chi_{[-\frac{a}{2\pi}, \frac{a}{2\pi}]}(x)$$

We will use  $a = 1, \frac{1}{3},$  and  $\frac{1}{5}$ . The relevant convolution was also computed above, but all we need is the fact that the support of

$$3\pi\chi_{[-\frac{1}{6\pi}, \frac{1}{6\pi}]} * 5\pi\chi_{[-\frac{1}{10\pi}, \frac{1}{10\pi}]}$$

is inside the interval  $[-\frac{1}{6\pi} - \frac{1}{10\pi}, \frac{1}{6\pi} + \frac{1}{10\pi}]$ . Thus, the integral of three *sinc* functions is equal to

$$\begin{aligned} \int_{\mathbb{R}} \pi\chi_{[-\frac{1}{2\pi}, \frac{1}{2\pi}]}(x) \cdot \left(3\pi\chi_{[-\frac{1}{6\pi}, \frac{1}{6\pi}]} * 5\pi\chi_{[-\frac{1}{10\pi}, \frac{1}{10\pi}]}\right)(x) dx &= \pi \cdot 3\pi \cdot 5\pi \int_{-1/\pi}^{1/\pi} \left(\chi_{[-\frac{1}{6\pi}, \frac{1}{6\pi}]} * \chi_{[-\frac{1}{10\pi}, \frac{1}{10\pi}]}\right)(x) dx \\ &= \pi \cdot 3\pi \cdot 5\pi \int_{\mathbb{R}} \left(\chi_{[-\frac{1}{6\pi}, \frac{1}{6\pi}]} * \chi_{[-\frac{1}{10\pi}, \frac{1}{10\pi}]}\right)(x) dx \end{aligned}$$

since  $[-1/2\pi, 1/2\pi]$  contains the support of the convolution. Observing that (invoking Fubini-Tonelli as necessary),

$$\int_{\mathbb{R}} (f * g)(x) dx = \int_{\mathbb{R}} \int_{\mathbb{R}} f(x-y)g(y) dx dy = \int_{\mathbb{R}} \int_{\mathbb{R}} f(x)g(y) dx dy = \int_{\mathbb{R}} f(x) dx \cdot \int_{\mathbb{R}} g(x) dy$$

the integral of the convolution is

$$\int_{\mathbb{R}} \chi_{[-\frac{1}{6\pi}, \frac{1}{6\pi}]} \cdot \int_{\mathbb{R}} \chi_{[-\frac{1}{10\pi}, \frac{1}{10\pi}]} = \frac{1}{3\pi} \cdot \frac{1}{5\pi}$$

Thus, the whole is

$$\pi \cdot 3\pi \cdot 5\pi \cdot \frac{1}{3\pi} \cdot \frac{1}{5\pi} = \pi$$

Similarly, the integral of  $f_1 * \dots * f_n$  is the product of the integrals  $\int f_i$ . With the support of  $f_i$  inside  $[-a_i, a_i]$ , the support of the convolution is inside  $[-a_1 - \dots - a_n, a_1 + \dots + a_n]$ . Thus, since  $\frac{1}{3} + \frac{1}{5} + \dots + \frac{1}{13} < 1$ , the same argument shows that

$$\int_{\mathbb{R}} \frac{\sin x}{x} \cdot \frac{\sin x/3}{x/3} \cdot \frac{\sin x/5}{x/5} dx = \pi \quad (\text{for } 2n+1 = 3, 5, 7, 9, 11, 13)$$

but for  $2n+1 = 15$ , the support of the Fourier transform of  $\frac{\sin x}{x}$  no longer contains the support of the convolution. ///

[04.6] Compute  $e^{-\pi x^2} * e^{-\pi x^2}$  and  $\frac{\sin x}{x} * \frac{\sin x}{x}$ . (Be careful what you assert:  $\frac{\sin x}{x}$  is not in  $L^1(\mathbb{R})$ .)

**Discussion:** The idea is to invoke  $f * g = (\widehat{f \cdot g})^\wedge$  for *even* functions  $f, g \in L^1$ , since for even functions the inverse Fourier transform is the same as the forward Fourier transform. Conveniently, Gaussians are in  $L^1 \cap L^2$ , and, from above, have Fourier transforms which are again Gaussians:

$$e^{-\widehat{\pi a x^2}}(\xi) = \frac{1}{\sqrt{a}} e^{-\pi \xi^2/a} \quad (\text{for } a > 0)$$

so

$$e^{-\pi x^2} * e^{-\pi x^2}(\xi) = e^{-\pi x^2} \cdot \widehat{e^{-\pi x^2}}(\xi) = e^{-2\pi x^2}(\xi) = \frac{1}{\sqrt{2}} e^{-\pi \xi^2/2}$$

For the other example, the bound  $|f * g|_{L^1} \leq |f|_{L^p} \cdot |g|_{L^q}$  for conjugate exponents  $p, q$  shows that  $f * g \in L^1$  for  $f, g \in L^2$ . Thus, the same identity holds for  $f, g \in L^2$ , with the Plancherel extension of Fourier transform. That is,  $\widehat{f}$  and  $\widehat{g}$  need not be the literal integrals for the Fourier transform, but its extension by continuity to  $L^2$ . Above, we computed the Fourier transform of characteristic functions of intervals:

$$\chi_{[-a, a]} \widehat{a}(\xi) = \frac{\sin 2\pi a \xi}{\pi \xi}$$

Thus,

$$(\pi \cdot \chi_{[-1/2\pi, 1/2\pi]}) \widehat{(\xi)} = \frac{\sin \xi}{\xi}$$

Then

$$\begin{aligned} \left(\frac{\sin x}{x} * \frac{\sin x}{x}\right)(\xi) &= \left((\pi \cdot \chi_{[-1/2\pi, 1/2\pi]}) \cdot (\pi \cdot \chi_{[-1/2\pi, 1/2\pi]})\right) \widehat{(\xi)} \\ &= \pi \cdot (\pi \cdot \chi_{[-1/2\pi, 1/2\pi]}) \widehat{(\xi)} = \pi \cdot \frac{\sin \xi}{\xi} \end{aligned}$$

///

[04.7] Prove that every  $f \in C_c^o(\mathbb{R})$  can be uniformly approximated (in sup norm) arbitrarily well as superpositions of Gaussians: given  $\varepsilon > 0$ , there is  $g \in C_c^o(\mathbb{R})$  and sufficiently large  $n$  such that

$$\sup_{x \in \mathbb{R}} \left| f(x) - \int_{\mathbb{R}} g(\xi) \cdot n e^{-\pi n^2 (\xi - x)^2} d\xi \right| < \varepsilon$$

**Discussion:** This is an instance of an *approximate identity* and the basic property of such. Namely, for an approximate identity  $\{\varphi_n\}$  on  $\mathbb{R}$  and  $f \in C_c^o(\mathbb{R})$ , we have

$$\sup_{x \in \mathbb{R}} \left| f(x) - \int_{\mathbb{R}} \varphi_n(\xi) \cdot f(x + \xi) d\xi \right| \longrightarrow 0 \quad (\text{as } n \rightarrow +\infty)$$

By replacing  $\xi$  by  $\xi - x$  in the integral, we have

$$\sup_{x \in \mathbb{R}} \left| f(x) - \int_{\mathbb{R}} f(\xi) \cdot \varphi_n(\xi - x) d\xi \right| \longrightarrow 0 \quad (\text{as } n \rightarrow +\infty)$$

Rather than reproving this general assertion in the example at hand, we simply clarify the interpretation in terms of approximate identities. That is, with  $\varphi_1(x) = e^{-\pi x^2}$ , we that the sequence  $\varphi_n(x) = n \cdot \varphi_1(n \cdot x)$  is an approximate identity. More generally, we prove

[0.1] **Claim:** Let  $\varphi \in C^o(\mathbb{R})$  be a non-negative  $\mathbb{R}$ -valued function, with  $\int_{\mathbb{R}} \varphi = 1$ . Then  $\varphi_n(x) = n \cdot \varphi(n \cdot x)$  is an approximate identity.

*Proof:* The non-negative real-valued-ness is of course immediate. The integral of  $\varphi_n$  is

$$\int_{\mathbb{R}} \varphi_n(x) dx = \int_{\mathbb{R}} n \cdot \varphi(n \cdot x) dx = \int_{\mathbb{R}} n \cdot \varphi(x) \frac{dx}{n} = \int_{\mathbb{R}} \varphi(x) dx = 1$$

by replacing  $x$  by  $x/n$  in the integral. Finally, to see that the masses of the  $\varphi_n$  bunch up near 0: Since  $\varphi \geq 0$  and

$$\lim_n \int_{-\sqrt{n}}^{\sqrt{n}} \varphi(x) dx = \int_{\mathbb{R}} \varphi(x) dx = 1$$

given  $\varepsilon > 0$  there is sufficiently large  $n_o$  such that for all  $n \geq n_o$

$$1 \leq \lim_n \int_{-\sqrt{n}}^{\sqrt{n}} \varphi(x) dx > 1 - \varepsilon$$

Then, by replacing  $x$  by  $x/n$  in the integral,

$$\int_{-\frac{1}{\sqrt{n}}}^{\frac{1}{\sqrt{n}}} \varphi_n(x) dx = \int_{-\frac{1}{\sqrt{n}}}^{\frac{1}{\sqrt{n}}} n \cdot \varphi(n \cdot x) dx = \int_{-\sqrt{n}}^{\sqrt{n}} \varphi(x) dx > 1 - \varepsilon$$

The verifies the bunching-up property. ///

[04.8] Without worrying too much about identifying the finite, positive constant  $\int_{\mathbb{R}} \frac{(\sin x)^2}{x^2} dx$ , prove that, for given  $f \in C_c^o(\mathbb{R})$ , given  $\varepsilon > 0$ , there is sufficiently large  $n$  and a function  $g \in C_c^o(\mathbb{R})$  such that

$$\sup_{x \in \mathbb{R}} \left| f(x) - \int_{\mathbb{R}} g(\xi) \cdot \frac{(\sin n(x - \xi))^2}{(x - \xi)^2} d\xi \right| < \varepsilon$$

**Discussion:** After the more general discussion of the previous example, this is just another such. ///

[04.9] Show that the *principal value* functional

$$f \longrightarrow PV \int_{\mathbb{R}} \frac{f(x)}{x} dx = \lim_{\varepsilon \rightarrow 0} \left( \int_{-\infty}^{-\varepsilon} \frac{f(x)}{x} dx + \int_{\varepsilon}^{\infty} \frac{f(x)}{x} dx \right)$$

is equal to

$$- \int_{\mathbb{R}} f'(x) \cdot \log |x| dx$$

for  $f$  *continuously* differentiable on  $\mathbb{R}$ , with hypotheses on the decay of  $f$  and  $f'$  at infinity.

**Discussion:** This is an exercise in careful integration by parts, in the course of which we discover reasonable hypotheses on  $f$  and  $f'$  so that the natural heuristic is a proof.

For fixed small  $\varepsilon > 0$  and large  $M > 0$ , integration by parts gives

$$\begin{aligned} & \int_{-M}^{-\varepsilon} \frac{f(x)}{x} dx + \int_{\varepsilon}^M \frac{f(x)}{x} dx \\ &= \left[ f(x) \cdot \log |x| \right]_{-M}^{-\varepsilon} + \left[ f(x) \cdot \log x \right]_{\varepsilon}^M - \int_{-M}^{-\varepsilon} f'(x) \cdot \log |x| dx - \int_{\varepsilon}^M f'(x) \cdot \log x dx \end{aligned}$$

The simplest way to make the boundary terms near  $\pm\infty$  vanish is that they vanish *individually*. For example, it does not suffice that  $f \in L^1(\mathbb{R})$  or  $L^2(\mathbb{R})$ , because such a hypothesis by itself does not assure that  $f(x) \cdot \log |x|$  goes to 0 at  $\pm\infty$ , since  $f$  could have narrower-and-narrower spikes parading out to infinity. It is true that an additional condition on the derivative might promise this asymptotic behavior of  $f$ , but let's not be toooo clever. So, exactly require that  $f(x) \cdot \log |x|$  goes to 0 at  $\pm\infty$ . In contrast, making

$$\lim_{\varepsilon \rightarrow 0^+} \left( f(\varepsilon) \cdot \log |\varepsilon| - f(-\varepsilon) \cdot \log |\varepsilon| \right) = 0$$

is the most subtle issue here. It cannot reasonably accomplished by having the individuals go to 0, unless we require some sort of vanishing of  $f$  at 0, which would be undesirable here. Here is where *differentiability* of

$f$  at 0 can be used: a Taylor-Maclaurin expansion with remainder ensures that  $f(\varepsilon) - f(-\varepsilon) = O(\varepsilon)$ . Since  $\lim_{\varepsilon \rightarrow 0^+} \varepsilon \log \varepsilon = 0$ , the combination of these boundary terms does go to zero.

For the limit as  $M \rightarrow +\infty$  of the individual integrals of  $f'(x) \cdot \log |x|$  to exist in a simple fashion, it suffices that  $\lim_{M \rightarrow \infty} \int_M^\infty f'(x) \cdot \log |x| dx = 0$  and  $\lim_{M \rightarrow \infty} \int_{-\infty}^{-M} f'(x) \cdot \log |x| dx = 0$ . The last question is about what it takes to make

$$\lim_{\varepsilon \rightarrow 0^+} \left( \int_\varepsilon^\infty f'(x) \cdot \log |x| dx + \int_{-\infty}^{-\varepsilon} f'(x) \cdot \log |x| dx \right) = \int_{\mathbb{R}} f'(x) \cdot \log |x| dx$$

Since  $\log |x|$  is locally integrable, for  $f'$  being merely *essentially bounded* on some interval  $[-\varepsilon_o, \varepsilon_o]$ , e.g., *continuous*, the two individual integrals  $\int_\varepsilon^\varepsilon f'(x) \cdot \log |x| dx$  and  $\int_{-\varepsilon}^{-\varepsilon'} f'(x) \cdot \log |x| dx$  for  $0 < \varepsilon' < \varepsilon$  go to zero as  $\varepsilon \rightarrow 0^+$ . Thus, with these natural sufficient constraints, we have the indicated identity. ///

[04.10] Let  $\psi_n(x) = e^{2\pi i n x}$ . Let  $\delta_{\mathbb{Z}}$  be the *Dirac comb*, that is, a periodic version of Dirac's  $\delta$ , describable as having Fourier series

$$\delta_{\mathbb{Z}} = \sum_{n \in \mathbb{Z}} 1 \cdot \psi_n \quad (\text{converging in } H^{-1}(\mathbb{T}) \text{ or even } H^{-\frac{1}{2}-\varepsilon}(\mathbb{T}) \text{ for all } \varepsilon > 0)$$

With  $\lambda \notin \mathbb{R}$ , show that the differential equation

$$u'' - \lambda \cdot u = \delta_{\mathbb{Z}}$$

has a periodic solution  $u \in H^{\frac{3}{2}-\varepsilon}(\mathbb{T}) \subset C^0(\mathbb{T})$ , using Fourier series, *by division*. Show that the equation  $v'' - \lambda v = f$  is solved by

$$v(x) = \int_{\mathbb{T}} u(x-t) f(t) dt = \int_0^1 u(x-t) f(t) dt$$

**Discussion:** Using the *spectral* characterization of the  $H^{-\frac{1}{2}-\varepsilon}(\mathbb{T})$  norm,

$$\left\| \sum_{n \in \mathbb{Z}} 1 \cdot \psi_n \right\|_{H^{-\frac{1}{2}-\varepsilon}}^2 = \sum_{n \in \mathbb{Z}} |1|^2 \cdot (1+n^2)^{-\frac{1}{2}-\varepsilon}$$

which is convergent for all  $\varepsilon > 0$ , by comparison to  $\sum_{n \neq 0} 1/n^2$ . So that Fourier series converges in  $H^{-\frac{1}{2}-\varepsilon}(\mathbb{T})$  and produces a *generalized function* there.

The extension by continuity of  $d/dx$  from  $C^\infty(\mathbb{T}) \rightarrow C^\infty(\mathbb{T})$  to  $\frac{d}{dx} : H^s(\mathbb{T}) \rightarrow H^{s-1}(\mathbb{T})$  is continuous, by design. Similarly,  $\frac{d^2}{dx^2} : H^s(\mathbb{T}) \rightarrow H^{s-2}(\mathbb{T})$  is continuous. That is, since infinite sums are the corresponding limits of finite partial sums, this continuity means that termwise differentiation is correct. Let  $u = \sum_n c_n \psi_n$ , and solve, dropping the tilde from the notation,

$$\sum_{n \in \mathbb{Z}} 1 \cdot \psi_n = \delta_{\mathbb{Z}} = u'' - \lambda u = \sum_n c_n \left( \frac{d^2}{dx^2} - \lambda \right) \psi_n = \sum_n c_n (-4\pi^2 n^2 - \lambda) \cdot \psi_n$$

Equating coefficients,  $c_n = 1/(-4\pi^2 n^2 - \lambda)$ , for  $\lambda$  not equal to  $-4\pi^2 n^2$  for integer  $n$ . Another easy estimate shows that this  $u$  has gained 2 Sobolev indices, so is in  $H^{\frac{3}{2}-\varepsilon}(\mathbb{T})$ .

By Sobolev imbedding/inequality,  $H^s(\mathbb{T}) \subset C^0(\mathbb{T})$  for all  $s > \frac{1}{2}$ , so the solution is continuous (and, in fact, satisfies a further Lipschitz condition).

To see that  $v'' - \lambda v = f$  is solved by

$$v(x) = \int_{\mathbb{T}} u(x-t) f(t) dt = \int_0^1 u(x-t) f(t) dt$$

take  $f$  such that  $\widehat{f} \in \ell_{\mathbb{Z}}^1(\mathbb{T})$ , meaning that  $\sum_n |\widehat{f}(n)| < \infty$ . A somewhat stronger, more intuitive assumption is that  $f \in C^2(\mathbb{T})$ , and then by integration by parts

$$\widehat{f''}(n) = \int_{\mathbb{T}} e^{-2\pi i n x} f''(x) dx = \int_{\mathbb{T}} (-2\pi i n)^2 e^{-2\pi i n x} f(x) dx = (2\pi i n)^2 \cdot \widehat{f}(n)$$

(On the circle  $\mathbb{T}$ , and/or for  $\mathbb{Z}$ -periodic functions, there are no boundary terms in integration by parts.) We do not even to invoke Riemann-Lebesgue, since  $|\widehat{f''}(n)|$  is *bounded*, so there is a constant  $C$  such that  $|\widehat{f}(n)| \leq C/n^2$ , so  $\widehat{f} \in \ell_{\mathbb{Z}}^1$ .

Then Fubini-Tonelli assures the legitimacy of interchanging sum and limit: [1]

$$\begin{aligned} \int_{\mathbb{T}} u(x-t) f(t) dt &= \int_{\mathbb{T}} \sum_m \frac{1}{-4\pi^2 m^2 - \lambda} e^{2\pi i m(x-t)} \cdot \sum_n \widehat{f}(n) e^{2\pi i n t} dt \\ &= \sum_{m,n} \frac{\widehat{f}(n)}{-4\pi^2 m^2 - \lambda} \int_{\mathbb{T}} e^{2\pi i m(x-t)} e^{2\pi i n t} dt = \sum_n \frac{\widehat{f}(n)}{-4\pi^2 n^2 - \lambda} e^{2\pi i n x} \end{aligned}$$

by mutual orthogonality of distinct exponentials (in every Sobolev space). By Riemann-Lebesgue,  $\widehat{f}(n) \rightarrow 0$ , so

$$\sum_n |\widehat{f}(n)|^2 \cdot (1+n^2)^s < \infty \quad (\text{for any } s < -\frac{1}{2})$$

so  $f \in H^{-1}(\mathbb{T})$ , for example. Application of the (extension of)  $\frac{d^2}{dx^2} - \lambda$  *termwise* (again, justified by continuity of the extension) produces the Fourier expansion of  $f$ . ///

[0.2] **Remark:** The latter example illustrates the utility of using *generalized functions* even in a discussion that seems not to refer to them: there was no need to *guess* the function  $u(x-t)$  (sometimes called a *Green's function*) solving the differential equation, since we *solved* for it using the Fourier expansion of  $\delta_{\mathbb{Z}}$  that only converges in  $H^{-\frac{1}{2}-\varepsilon}(\mathbb{T})$ .

[1] In fact, we will see later that for  $u$  continuous and  $f$  in any Sobolev space, the interchange is justified.