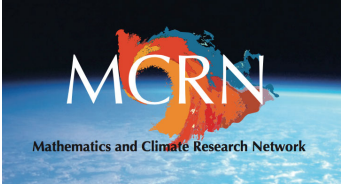



Budyko's Model and Snowball Earth

Richard McGehee
School of Mathematics
University of Minnesota
Mathematics of Climate Seminar
September 17, 2024



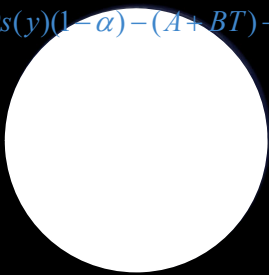
<https://sites.google.com/view/math-climate>

Budyko's Model

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T} - T)$$


Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball


$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T} - T)$$


Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Is it possible for Earth to become completely covered in ice? (Snowball Earth)

Did it ever happen?



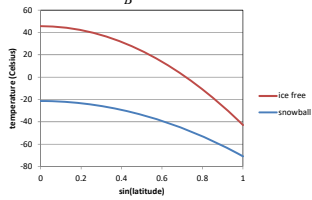
Mathematics of Climate Seminar 9/17/2024

Budyko's Model

Latitude Dependence

$$R \frac{\partial T(y,t)}{\partial t} = Qs(y)(1-\alpha) - (A+BT(y,t))$$

Equilibrium Solution

$$T^*(y) = \frac{Qs(y)(1-\alpha) - A}{B}$$


Tung*
 $\alpha = 0.32$: ice free
 $\alpha = 0.62$: snowball

Note that $\alpha = 0.32$ is in the range of current Earth.

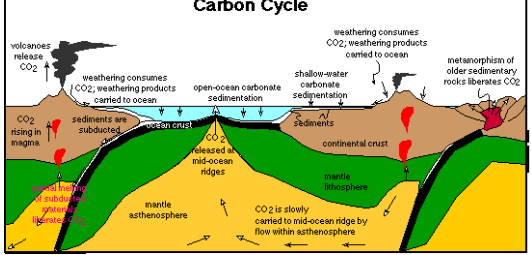
* K.K. Tung, *Topics in Mathematical Modeling*, Princeton U. Press, 2007

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Earth's Carbon Cycle

Long-Term Carbon Cycle



http://www.carleton.edu/departments/geol/DaveSTELLA/Carbon/long_term_carbon.htm

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Earth's Carbon Cycle

Volcanos emit CO₂.

Silicate weathering carries carbon to the ocean.



Carbon sinks to the bottom of the ocean by the "biological pump" and by the precipitation of calcium carbonate.

The carbon is captured in the sediment.

The sediment is subducted beneath the continental crust by plate tectonics.

Volcanic activity released carbon from the carbonate rocks in the form of CO₂.

Repeat.

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball



Earth's Carbon Cycle

Silicate Weathering

Rainwater containing dissolved CO₂ falling on silicate rocks replaces a silicon atom with a carbon atom, ultimately producing calcium carbonate (limestone) and silicon dioxide (quartz). For example, calcium silicate (Wollastonite):

$$CaSiO_3 + CO_2 \rightarrow CaCO_3 + SiO_2$$

Under volcanic conditions, the carbon atom is replaced by a silicon atom, completing the long term carbon cycle.

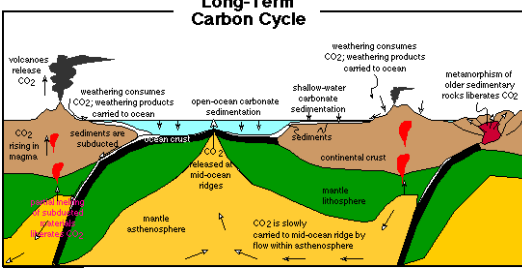
$$CaCO_3 + SiO_2 \rightarrow CaSiO_3 + CO_2$$



Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Earth's Carbon Cycle

Long-Term Carbon Cycle

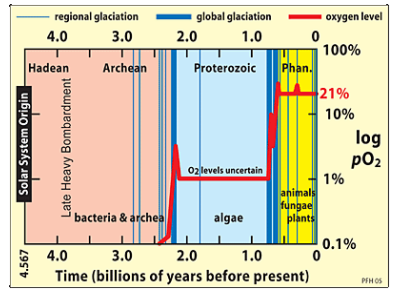


http://www.carleton.edu/departments/geol/DaveSTELLA/Carbon/long_term_carbon.htm



Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball




There is evidence that Snowball Earth has occurred, the last time about 600 million years ago.

<http://www.snowballearth.org/when.html>



Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball




The continents were clustered near the equator.

Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75



Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball




"Ice-rafted debris" occurred in ocean sediments near the equator, indicating large equatorial glaciers calving icebergs.

Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75





Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball




Large limestone deposits "cap carbonates" are found immediately above the glacial debris, indicating a rapid warming period following the snowball.



Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball



If Earth ever was a snowball, how did we escape?

Idea:


When the Earth is ice-covered, silicate weathering almost disappears, but volcanic activity continues, allowing for a build-up of CO₂ in the atmosphere.

If the ice starts to melt, the ocean opens up and starts to reabsorb the CO₂, which warms the planet, causing more ice to melt, causing more CO₂ to be absorbed in the ocean, etc. Rapid positive feedback leads to an ice-free Earth.

Mathematics of Climate Seminar 9/17/2024

Budyko's Model

Budyko's Equation



M. I. Budyko, "The effect of solar radiation variations on the climate of the Earth," *Tellus XXI*, 611-619, 1969.

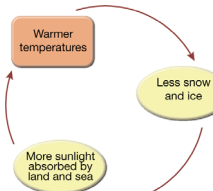
$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T}-T)$$

Labels for the equation: R (heat capacity), $\frac{\partial T}{\partial t}$ (surface temperature), $Qs(y)$ (insolation), $(1-\alpha)$ (albedo), $(A+BT)$ (OLR), $C(\bar{T}-T)$ (heat transport). $\bar{T} = \int_0^1 T(y) dy$ (sin(latitude)).

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice-albedo Feedback



temperature warms
ice melts
albedo decreases
more sunlight absorbed
temperature warms
REPEAT

Why would it stop?

Where is this feedback in Budko's equation?

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T}-T)$$

α shouldn't be a constant.

<http://www.i-fink.com/melting-polar-ice/>

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Budyko's Equation

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T}-T)$$

Where is the ice-albedo feedback?

Ice-albedo Feedback

albedo of ice: 0.62
albedo of land and water: 0.32

Assumption: there is a single boundary between ice and no ice occurring at $y = \eta$.

$$\alpha = \alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Budyko's Equation

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T}-T)$$

Where is the ice-albedo feedback?

Ice-albedo Feedback


albedo of ice: 0.62
albedo of land and water: 0.32

Assumption: there is a single boundary between ice and no ice occurring at $y = \eta$.

$$\alpha = \alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$$

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y, \eta)) - (A+BT) + C(\bar{T}-T)$$

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Ice-albedo Feedback

albedo of ice: 0.62
albedo of land and water: 0.32


Assumption: there is a single boundary between ice and no ice occurring at $y = \eta$.

$$\alpha = \alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$$

What determines the ice line η ?
Is it a dynamic variable or a parameter?
For now, we consider it to be a parameter.

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Budyko's Equilibrium

$$R \frac{\partial T}{\partial t}(y, t) = Qs(y)(1 - \alpha(y, \eta)) - (A + BT(y, t)) + C(\bar{T}(t) - T(y, t))$$

Equilibrium:
 $Qs(y)(1 - \alpha(y, \eta)) - (A + BT^*(y)) + C(\bar{T}^* - T^*(y)) = 0$

Assume we know the global mean temperature.

Solve for $T^*(y)$


$$Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^* - (BT^*(y)) + C(-T^*(y)) = 0$$

$$(B + C)T^*(y) = Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*$$

$$T^*(y) = \frac{Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*}{B + C}$$

Now we can compute the global mean temperature!

Mathematics of Climate Seminar 9/17/2024




Budyko's Snowball

Budyko's Equilibrium *We need to compute this.*

$$T^*(y) = \frac{Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*}{B + C}$$

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Budyko's Equilibrium *We need to compute this.*

$$T^*(y) = \frac{Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*}{B + C}$$

$$(B + C)T^*(y) = Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*$$

Integrate:


$$\int_0^1 (B + C)T^*(y) dy = \int_0^1 (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*) dy$$

$$(B + C) \int_0^1 T^*(y) dy = Q \left(\int_0^1 s(y) dy - \int_0^1 \alpha(y, \eta) s(y) dy \right) - A \int_0^1 dy + C\bar{T}^* \int_0^1 dy$$

$$\bar{T}^* = \frac{Q(1 - \bar{\alpha}(\eta)) - A}{B}$$

Success!

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Budyko's Equilibrium

$$T^*(y) = \frac{Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*}{B + C}$$

$$(B + C)T^*(y) = Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*$$

Integrate:

$$\int_0^1 (B + C)T^*(y) dy = \int_0^1 (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*) dy$$


$$(B + C) \int_0^1 T^*(y) dy = Q \left(\int_0^1 s(y) dy - \int_0^1 \alpha(y, \eta) s(y) dy \right) - A \int_0^1 dy + C\bar{T}^* \int_0^1 dy$$

$$\bar{T}^* = \frac{Q(1 - \bar{\alpha}(\eta)) - A}{B}$$

Note the dependence on the parameter η .

$$T^*(y, \eta) = \frac{Qs(y)(1 - \bar{\alpha}(\eta)) - A + C\bar{T}^*}{B + C}, \text{ where } \bar{T}^*(\eta) = \frac{Q(1 - \bar{\alpha}(\eta)) - A}{B}$$

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Budyko's Equilibrium

$$T^*(y, \eta) = \frac{Qs(y)(1 - \bar{\alpha}(\eta)) - A + C\bar{T}^*}{B + C}, \text{ where } \bar{T}^*(\eta) = \frac{Q(1 - \bar{\alpha}(\eta)) - A}{B}$$

Can we compute the global albedo?


recall: $\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$

global albedo $\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta) s(y) dy = \int_0^\eta \alpha_1 s(y) dy + \int_\eta^1 \alpha_2 s(y) dy$
 $= \alpha_1 S(\eta) + \alpha_2 (1 - S(\eta))$

where $S(\eta) = \int_0^\eta s(y) dy = \int_0^\eta (1 - 0.241(3y^2 - 1)) dy = \eta - 0.241(\eta^3 - \eta)$
Chylek & Coakley

Time to summarize.

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Budyko's Equilibrium

$$R \frac{\partial T}{\partial t}(y, t) = Qs(y)(1 - \alpha(y, \eta)) - (A + BT(y, t)) + C(\bar{T}(t) - T(y, t))$$

Equilibrium temperature distribution:


$$T^*(y, \eta) = \frac{Qs(y)(1 - \bar{\alpha}(\eta)) - A + C\bar{T}^*(\eta)}{B + C}$$

$$\bar{T}^*(\eta) = \frac{Q(1 - \bar{\alpha}(\eta)) - A}{B}$$


$$\bar{\alpha}(\eta) = \alpha_s S(\eta) + \alpha_i (1 - S(\eta))$$

$$s(y) = 1 - 0.241(3y^2 - 1)$$

$$S(\eta) = \eta - 0.241(\eta^2 - \eta)$$




Mathematics of Climate Seminar 9/17/2024




Budyko's Snowball

Budyko's Equilibria

For each fixed η , there is an equilibrium solution for Budyko's equation.

$$T_\eta^*(y) = \frac{1}{B + C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}_\eta^*)$$


Mathematics of Climate Seminar 9/17/2024

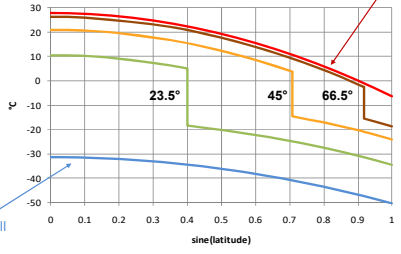



Budyko's Snowball


Budyko's Equilibria

$$T_\eta^*(y) = \frac{1}{B + C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}_\eta^*)$$

For each fixed η , there is an equilibrium solution for Budyko's equation.

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Stability of Equilibria

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

Let X be the space of functions where T lives. (e.g. $L^2([0,1])$)

Let

$$L: X \rightarrow X: LT = C\bar{T} - (B + C)T,$$


$$f(y) = Qs(y)(1 - \alpha(y)) - A$$

Budyko's equation can be written as a linear vector field on X .


$$R \frac{dT}{dt} = f + LT$$

The operator L has only point spectrum, with all eigenvalues negative. Therefore, *all solutions are stable*. True for any albedo function.

experts only



Mathematics of Climate Seminar 9/17/2024



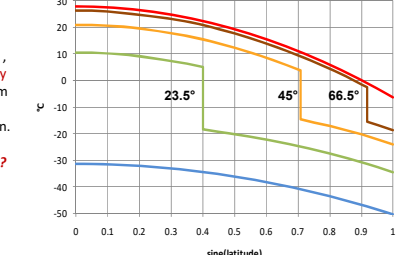

Budyko's Snowball

Stability of Equilibria


$$R \frac{\partial T}{\partial t}(y, t) = Qs(y)(1 - \alpha(y, \eta)) - (A + BT(y, t)) + C(\bar{T}(t) - T(y, t))$$

For each fixed η , there is a **globally stable** equilibrium solution for Budyko's equation.

How to pick one?

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Ice Albedo Feedback


Idea

If we artificially hold the ice line at a fixed latitude, then the surface temperature will come to an equilibrium.

However,

if the temperature at the ice line is high, we would expect ice to melt and the ice line to retreat to higher latitudes. If the temperature is low, we would expect ice to form and the ice line to advance to lower latitudes.

How to model this expectation?



Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice Albedo Feedback

For each fixed η , there is a stable equilibrium solution for Budyko's equation.

Standard assumption: Permanent ice forms if the annual average temperature is below $T_c = -10^\circ\text{C}$ and melts if the annual average temperature is above T_c .

Additional condition: The average temperature across the ice boundary must be the critical temperature T_c .

$$\frac{1}{2}(T_\eta^*(\eta+) + T_\eta^*(\eta-)) = T_c = -10$$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice Albedo Feedback

ice line condition: $\frac{1}{2}(T_\eta^*(\eta+) + T_\eta^*(\eta-)) = T_c = -10$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(T - T_c)$$

Equilibrium: $T_\eta^*(y) = \frac{1}{B+C}(Qs(y)(1 - \alpha(y, \eta)) - A + CT_\eta^*)$

Ice line condition: $\frac{1}{2}(T_\eta^*(\eta+) + T_\eta^*(\eta-)) = T_c = -10$

Albedo: $\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases} \quad \alpha(\eta-, \eta) = \alpha_1, \quad \alpha(\eta+, \eta) = \alpha_2$

$$T_\eta^*(\eta+) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_2) - A + CT_\eta^*) \quad T_\eta^*(\eta-) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_1) - A + CT_\eta^*)$$

Ice line condition: $\frac{1}{2}(T_\eta^*(\eta+) + T_\eta^*(\eta-)) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_0) - A + CT_\eta^*) = T_c = -10$

where: $\alpha_0 = \frac{\alpha_1 + \alpha_2}{2} = 0.47$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(T - T_c)$$

Ice line condition: $\frac{1}{B+C}(Qs(\eta)(1 - \alpha_0) - A + CT_\eta^*) = T_c = -10$

Rewrite: $h(\eta) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_0) - A + CT_\eta^*) - T_c = 0$

Recall equilibrium GMT: $\bar{T}_\eta = \frac{1}{B}(Q(1 - \bar{\alpha}(\eta)) - A)$

Recall average albedo: $\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta) s(y) dy = \alpha_2 - (\alpha_2 - \alpha_1)S(\eta) = 0.62 - 0.3S(\eta)$

where: $S(\eta) = \int_0^\eta s(y) dy = \eta - 0.241(\eta^5 - \eta)$

$$h(\eta) = \frac{Q}{B+C} \left(s(\eta)(1 - \alpha_0) + \frac{C}{B}(1 - \alpha_2 + (\alpha_2 - \alpha_1)S(\eta)) \right) - \frac{A}{B} - T_c = 0$$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(T - T_c)$$

The additional condition: $\frac{1}{2}(T_\eta^*(\eta+) + T_\eta^*(\eta-)) = T_c = -10$

can be written: $h(\eta) = \frac{Q}{B+C} \left(s(\eta)(1 - \alpha_0) + \frac{C}{B}(1 - \alpha_2 + (\alpha_2 - \alpha_1)S(\eta)) \right) - \frac{A}{B} - T_c = 0$

Two equilibria (zeros of h) satisfy the additional condition.

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Ice Albedo Feedback

Equilibrium temperature profiles $T_\eta^*(y) = \frac{1}{B+C}(Qs(y)(1 - \alpha(y, \eta)) - A + CT_\eta^*)$

Interesting Solutions:

- small cap
- large cap
- ice free
- snowball

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Dynamics of the Ice Line

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Idea:
If the average temperature across the ice line is above the critical temperature, some ice will melt, moving the ice line toward the pole. If it is below the critical temperature, the ice will advance toward the equator.

stationary ice melts stationary

Widiasih's equation: $\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Dynamics of the Ice Line

$$\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$$

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Widiasih's Theorem. For sufficiently small ε , the system has an attracting invariant curve given by the graph of a function $\Phi_\varepsilon : [0, 1] \rightarrow X$. On this curve, the dynamics are approximated by the equation

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

experts only

unstable stable

Esther R. Widiasih, Dynamics of the Budyko Energy Balance Model, SIAM J. Appl. Dyn. Syst., 12(4), 2068-2092.

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Budyko-Widiasih Model

Temperature profiles

Esther Widiasih

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Budyko-Widiasih Model

Esther Dick

Oahu, February 2020

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Summary

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

surface temperature sin(latitude) $\bar{T} = \int_0^1 T(y) dy$

heat capacity insolation albedo OLR heat transport

reduces to

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

$$\frac{d\eta}{dt} = \varepsilon h(\eta) = \varepsilon \left(\frac{Q}{B+C} (s(\eta)(1 - \alpha_0) + \frac{C}{B} (1 - \alpha_2 + (\alpha_2 - \alpha_1) S(\eta))) - \frac{A}{B} - T_c \right)$$

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Budyko-Widiasih Model

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

surface temperature sin(latitude) $\bar{T} = \int_0^1 T(y) dy$

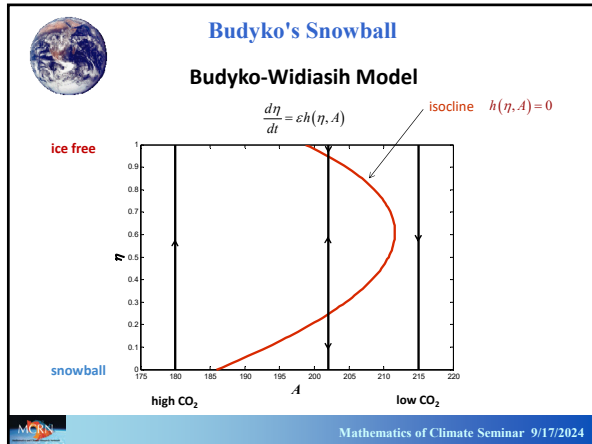
heat capacity insolation albedo OLR heat transport

What about the greenhouse effect?

$A + BT$ is the outgoing long wave radiation. This term decreases if the greenhouse gases increase. We view A as a parameter.

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$

Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Is it possible for Earth to become a snowball?
If so, how does it escape?

Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$

What if A is a dynamical variable?

Simple equation:

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c) \leftarrow \text{silicate weathering}$$

$0 < \tilde{\varepsilon} \ll \varepsilon \ll 1$

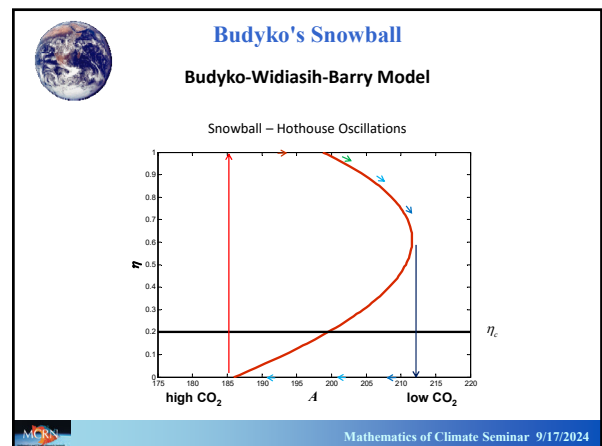
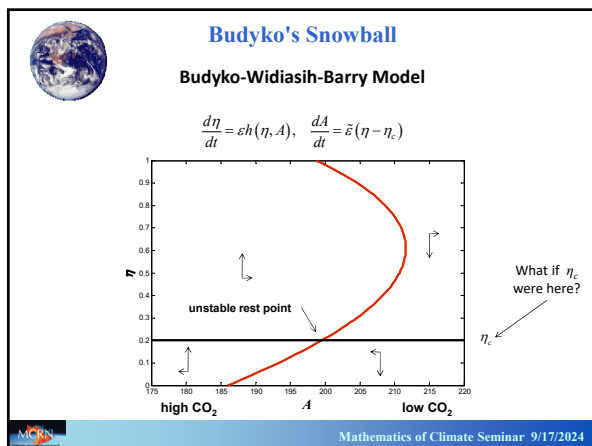
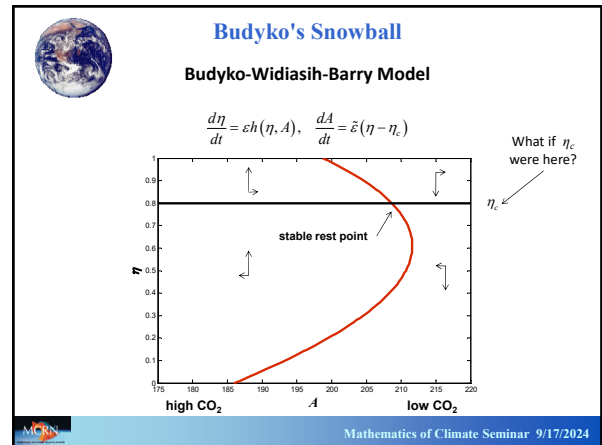
New system:

$$\begin{cases} \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c) \\ \frac{d\eta}{dt} = \varepsilon h(\eta, A) \end{cases}$$

Anna Barry
Postdoc 2012-14

Anna M. Barry, Esther Widiasih, & Richard McGehee, *Discrete & Continuous Dynamical Systems Series B* 22 (2017), 2447-2463.

Mathematics of Climate Seminar 9/17/2024





Budyko's Snowball


Banded Iron Formations



Banded iron deposits form when the ocean oscillates between oxygen-rich and oxygen-poor. The formation on the right is from northern Minnesota.

<https://www.usgs.gov/media/images/minnesota-banded-iron>

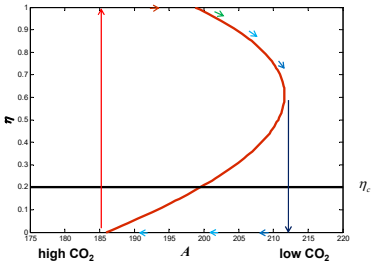
Mathematics of Climate Seminar 9/17/2024



Budyko's Snowball

Budyko-Widiasih-Barry Model

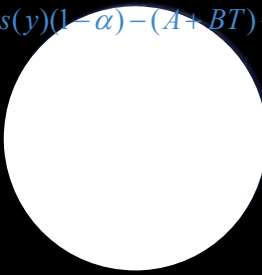
Snowball – Hothouse Oscillations



Mathematical Questions
Is there a theory about the dynamics of vector fields on manifolds with boundary?
Existence of solutions?
Uniqueness?
Bifurcation theory?

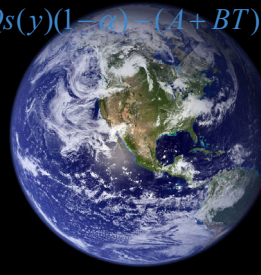
Mathematics of Climate Seminar 9/17/2024

Budyko's Snowball

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T} - T)$$


Mathematics of Climate Seminar 9/17/2024

Budyko's Model

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T} - T)$$


Mathematics of Climate Seminar 9/17/2024