

(116)

Two more important properties of $\omega \mapsto d\omega$:

THM 6.7.7: (i) For $\varphi \in A^k(\Omega)$ having coefficients in $C^2(\Omega)$, $d(d\varphi) = 0$.

6.7.8:

(ii) For $\varphi \in A^k(\Omega)$ and $\psi \in A^l(\Omega)$,

$$d(\varphi \wedge \psi) = d\varphi \wedge \psi + (-1)^k \varphi \wedge d\psi$$

proof: (ii) is EXER 6.7.11 on your HW!

For (i), it's enough to check it (by linearity of d) when

$$\varphi = f(x) dx_1 \wedge \dots \wedge dx_k$$

where one has

$$\begin{aligned} d(d\varphi) &= d\left(\sum_{i=1}^n \frac{\partial f}{\partial x_i} dx_1 \wedge \dots \wedge dx_i \wedge \dots \wedge dx_k\right) \\ &= \left(\sum_{i=1}^n \sum_{j=1}^n \frac{\partial^2 f}{\partial x_j \partial x_i} dx_j \wedge dx_i\right) \wedge dx_1 \wedge \dots \wedge dx_k \end{aligned}$$

$$\Rightarrow \text{since } dx_i \wedge dx_j = \begin{cases} 0 & \text{if } i=j \\ -dx_j \wedge dx_i & \text{if } i \neq j \end{cases}$$

$$\text{but } \frac{\partial^2 f}{\partial x_j \partial x_i} = \frac{\partial^2 f}{\partial x_i \partial x_j}$$

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EXAMPLE: $d(d x^2 y^3) = d(2x y^3 dx + 3x^2 y^2 dy) = 6x y^2 dy \wedge dx + 6x y^3 dx \wedge dy = 0$

§6.8 DV, grad, & curl

There are 3 operations in vector calculus of \mathbb{R}^3 and physics
that have nice interpretations/unifications via $\omega \mapsto d\omega$:

$$\text{For } f: \mathbb{R}^3 \rightarrow \mathbb{R}, \quad \text{grad } f := \begin{bmatrix} \frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y} \\ \frac{\partial f}{\partial z} \end{bmatrix} = " \nabla f " = \text{gradient of } f$$

$$\begin{aligned} \text{For a vector field } \vec{F}: \mathbb{R}^3 \rightarrow \mathbb{R}, \quad \text{curl } \vec{F} &:= " \nabla \times \vec{F} " = \begin{bmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{bmatrix} \times \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix} = \underline{\text{curl of } \vec{F}} \\ &= \begin{bmatrix} \cancel{\frac{\partial F_3}{\partial y}} - \frac{\partial F_2}{\partial z} \\ - \left(\cancel{\frac{\partial F_3}{\partial x}} - \frac{\partial F_1}{\partial z} \right) \\ \frac{\partial F_2}{\partial x} - \frac{\partial F_1}{\partial y} \end{bmatrix} \end{aligned}$$

$$\begin{aligned} \text{div } \vec{F} &:= " \nabla \cdot \vec{F} " = \begin{bmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{bmatrix} \cdot \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix} = \frac{\partial F_1}{\partial x} + \frac{\partial F_2}{\partial y} + \frac{\partial F_3}{\partial z} \\ &= \underline{\text{divergence of } \vec{F}} \end{aligned}$$

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THM 6.8.3: (i) For $f \in A^0(\mathbb{R}^3)$, i.e. $f: \mathbb{R}^3 \rightarrow \mathbb{R}$,

~~then~~ then $df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz = W_{\nabla f}$ = the work form of $\nabla f = \begin{bmatrix} \frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y} \\ \frac{\partial f}{\partial z} \end{bmatrix} \in A^1(\mathbb{R}^3)$

(ii) For $W_F = F_1 dx + F_2 dy + F_3 dz \in A^1(\mathbb{R}^3)$

$$= \text{Work form of } F = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix},$$

$$\text{then } dW_F = \left(\frac{\partial F_2}{\partial x} - \frac{\partial F_1}{\partial y} \right) dx \wedge dy + \left(\frac{\partial F_3}{\partial x} - \frac{\partial F_1}{\partial z} \right) dx \wedge dz + \left(\frac{\partial F_3}{\partial y} - \frac{\partial F_2}{\partial z} \right) dy \wedge dz$$

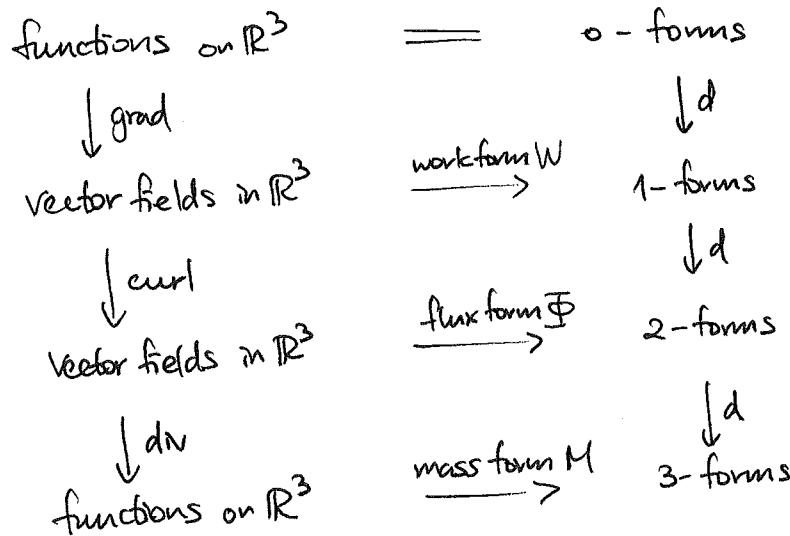
$$= \Phi_{\text{curl } F} = \text{the flux form of curl } F \in A^2(\mathbb{R}^3)$$

(iii) For $\Phi_F = F_1 dy \wedge dz - F_2 dx \wedge dz + F_3 dx \wedge dy \in A^2(\mathbb{R}^3)$

$$= \text{the flux 2-form of } F = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix},$$

$$\text{then } d\Phi_F = \left(\frac{\partial F_1}{\partial x} + \frac{\partial F_2}{\partial y} + \frac{\partial F_3}{\partial z} \right) dx \wedge dy \wedge dz = M_{\text{div } F}$$

$$= \text{the mass 3-form of div } F \in A^3(\mathbb{R}^3)$$

SUMMARY TABLE from p.630Note (RMK 6.8.6) that $d(d\varphi) = 0$ implies

$$\text{curl}(\text{grad } f) = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

"gradient flows (or conservative vector fields) are irrotational" $\uparrow \text{curl} = 0$

$F = \nabla f$ for some

$$\text{div}(\text{curl } \bar{F}) = 0$$

"a flow that comes from a curl is incompressible" $\uparrow \text{div} = 0$

(118) Why should a vector field $\bar{F} = \begin{pmatrix} F_1 \\ F_2 \\ F_3 \end{pmatrix}$ of the form $\bar{F} = \nabla f$ for some function $f: \mathbb{R}^3 \rightarrow \mathbb{R}$ be called conservative?

Dipping into §6.10, the Fundamental Thm. for Line Integrals says

THM 6.10.1: For an oriented curve $C \subset \mathbb{R}^2$ or \mathbb{R}^3 and $f \in C^1(U)$ in some open neighborhood U of C will have

$$\int_C df = f(b) - f(a) \quad (= \int_C f) \quad (\text{i.e. Stokes's Thm!})$$

$\int_C \bar{F} \cdot \bar{\gamma}'(t) dt$
 the line integral
 of \bar{F} over C

the answer only
 depends on the
 starting and endpoints b, a
 and the value of the potential
 function $f(x)$ there,
not the path $[C] = [\gamma(v)]$ taken to get there
conservative

What does $\text{curl } \bar{F}$ mean?

The book discusses the curl probe: put a paddle at $\bar{x} \in \mathbb{R}^3$ inside the flow with velocity vectors locally $\bar{F}(\bar{x}) = \begin{pmatrix} F_1 \\ F_2 \\ F_3 \end{pmatrix}$, and the locally torque exerted on the paddle will be proportional to $\text{curl } \bar{F}(\bar{x}) \cdot \bar{n}$ if the handle points toward \bar{n}



why? Call the paddle wheel blades \bar{v}_1, \bar{v}_2 , so that $\bar{n} = \bar{v}_1 \times \bar{v}_2$



$$\text{Then } \text{curl } \bar{F} \cdot \bar{n} = \text{curl } \bar{F}(\bar{x}) \cdot (\bar{v}_1 \times \bar{v}_2)$$

$$= \underbrace{\oint_{\text{curl } \bar{F}(\bar{x})} P_x(\bar{v}_1, \bar{v}_2)}_{= dW_F}$$

$$= \lim_{h \rightarrow 0} \frac{1}{h^2} \iint_{\partial P_x(h\bar{v}_1, h\bar{v}_2)} W_F$$



work done by \bar{F} as a force field going around the boundary of $P_F(h\bar{v}_1, h\bar{v}_2)$

