

# The Askey–Wilson Polynomials

## A Personal Journey

December 5, 2013

For a sequence of orthogonal polynomials you would like to find

1. Orthogonality measure.
2. Raising and lowering operators. Rodrigues formulas.
3. A second order Sturm–Liouville type equation.
4. Connection relations.
5. Coefficients in the linearization of products

$$p_m(x)p_n(x) = \sum_k c(m, n, k)p_k(x).$$

6. Generating functions  $\sum_n p_n(x) \lambda_n t^n$ , for suitable  $\lambda_n$ .
7. Poisson kernel  $\sum p_n(x) p_n(y) r^n$ ,  $p_n$  are orthonormal.
8. Asymptotics as the degree  $\rightarrow \infty$ .
9. Plancherel–Rotach asymptotics if applicable (unbounded support).

Notation:

$$(a; q)_n = \prod_{j=0}^{n-1} (1 - aq^j), \quad n = 0, 1, \dots, \text{ or } \infty,$$

$$(a_1, a_2, \dots, a_m; q)_n = \prod_{k=1}^m (a_k; q)_n$$

The Askey–Wilson memoir:

$$w(\cos \theta; \mathbf{a}) dx = \frac{(e^{2i\theta}, e^{-2i\theta}; q)_{\infty}}{\prod_{j=1}^4 (a_j e^{i\theta}, a_j e^{-i\theta}; q)_{\infty}} d\theta.$$

Notation:  $x = \cos \theta$  versus  $z = e^{i\theta}$ .

$$\bar{f}(z) = f((z + 1/z)/2).$$

For  $e(x) = x$ ,  $\bar{e}(z) = (z + 1/z)/2$ .

$$(\mathcal{D}_q f)(x) = \frac{\bar{f}(q^{1/2}z) - \bar{f}(q^{-1/2}z)}{(q^{1/2} - q^{-1/2})(z - 1/z)/2}.$$

For example

$$\mathcal{D}_q T_n(x) = \frac{q^{n/2} - q^{-n/2}}{q^{1/2} - q^{-1/2}} U_{n-1}(x).$$

I tried to understand the Askey–Wilson operator and the Askey–Wilson integral. I want to share some of the findings with you. Most of the work is joint work.

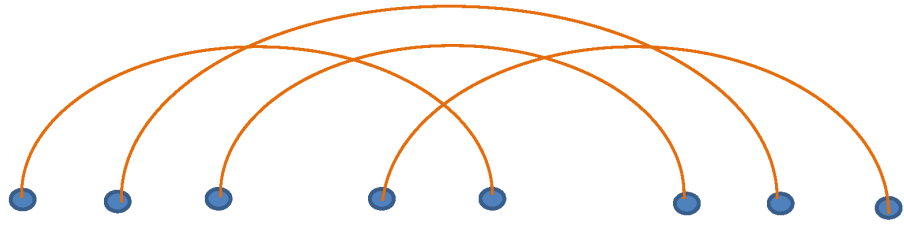
$$\int_0^\pi \frac{(e^{2i\theta}, e^{-2i\theta}; q)_\infty}{\prod_{j=1}^4 (a_j e^{i\theta}, a_j e^{-i\theta}; q)_\infty} d\theta = \frac{2\pi (a_1 a_2 a_3 a_4; q)_\infty}{(q; q)_\infty \prod_{1 \leq j < k \leq 4} (a_j a_k; q)_\infty}.$$

In 1893–1895 Rogers studied the  $q$ -Hermite polynomials

$$\sum_{n=0}^{\infty} H_n(\cos \theta | q) \frac{t^n}{(q; q)_n} = \frac{1}{(te^{i\theta}, te^{-i\theta}; q)_\infty}.$$

He proved the linearization formula

$$H_m(x|q)H_n(x|q) = \sum_{k=0}^{m \wedge n} \frac{(q; q)_m (q; q)_n}{(q; q)_k (q; q)_{m-k} (q; q)_{n-k}} H_{m+n-2k}(x|q)$$



Azor–Gillis–Victor and independently Godsil (1982) proved that the number of perfect matchings is

$$\frac{2^{-\sum_j n_j/2}}{\sqrt{\pi}} \int_{\mathbb{R}} e^{-x^2} \prod_{j=1}^m H_{n_j}(x) dx.$$

This means that the following integral is a  $q$ -analogue of the number of perfect matchings whose generating function is

$$\int_0^\pi \frac{(e^{2i\theta}, e^{-2i\theta}; q)_\infty}{\prod_{j=1}^m (a_j e^{i\theta}, a_j e^{-i\theta}; q)_\infty} d\theta.$$

the coefficient of  $\prod_{j=1}^m a_j^{n_j} / (q; q)_{n_j}$  is

$$\int_0^\pi (e^{2i\theta}, e^{-2i\theta}; q)_\infty \prod_{j=1}^m H_{n_j}(\cos \theta | q) d\theta.$$

After some scaling we see that this =

$$(1) \quad \sum_{\text{Perfect Matchings}} q^{\text{crossing number}}.$$

Ismail–Stanton–Viennot (1987).

**Problem:** The polynomial (1) is symmetric in  $n_1, n_2, \dots, n_m$ . It is

$$\int_{\mathbb{R}} w(x) \prod_{j=1}^m \tilde{H}_{n_j}(x|q) dx.$$

**Fact:**

$$\begin{aligned} & \frac{(q; q)_{\infty}}{2\pi} \int_0^{\pi} \frac{(e^{2i\theta}, e^{-2i\theta}; q)_{\infty}}{\prod_{j=1}^5 (a_j e^{i\theta}, a_j e^{-i\theta}; q)_{\infty}} d\theta \\ &= \frac{(a_1 a_2 a_3 a_5, a_2 a_3 a_4 a_5, a_1 a_4; q)_{\infty}}{\prod_{1 \leq j < k \leq 5} (a_j a_k; q)_{\infty}} \\ & \quad \times {}_3\phi_2 \left( \begin{matrix} a_2 a_3, a_2 a_5, a_3 a_5 \\ a_1 a_2 a_3 a_5, a_2 a_3 a_4 a_5 \end{matrix} \middle| q, a_1 a_4 \right). \end{aligned}$$

You get  ${}_3\phi_2$  transformations.



The Askey–Wilson operators: Null space is nontrivial, no two-sided inverse. What about one sided inverse.  $\frac{d}{dx} \int_a^x f(y)dy = f(x)$ . Analogue of  $\int_a^x$ . Also we need an analogue of integration by parts. This is essentially computing the adjoint of  $\mathcal{D}_q$ .

We shall use the inner product (Chebyshev)

$$\langle f, g \rangle := \int_{-1}^1 f(x) \overline{g(x)} \frac{dx}{\sqrt{1-x^2}}.$$

We require  $\bar{f}(z)$  to be defined for  $|q^{\pm 1/2}z| = 1$  as well as for  $|z| = 1$ . Thus  $\mathcal{D}_q$  is defined on:

$$H_\nu := \left\{ f : f((z + 1/z)/2) \text{ analytic in } q^\nu \leq |z| \leq q^{-\nu} \right\}.$$

**Theorem 1.** For  $f, g \in H_{1/2}$  we have

$$\begin{aligned} \langle \mathcal{D}_q f, g \rangle &= \frac{\pi \sqrt{q}}{1-q} [f((q^{1/2} + q^{-1/2})/2) \overline{g(1)} \\ &\quad - f(-(q^{1/2} + q^{-1/2})/2) \overline{g(-1)}] \\ &\quad - \left\langle f, \sqrt{1-x^2} \mathcal{D}_q(g(x)(1-x^2)^{-1/2}) \right\rangle. \end{aligned}$$

Analyze Sturm–Liouville equations. Brown–Evans–Ismail. Later work with Christiansen, Christainsen and Koelink, Brown–Chriatinsen.

One sided inverse:

$$\mathcal{D}_q T_n(x) = \frac{q^{n/2} - q^{-n/2}}{q^{1/2} - q^{-1/2}} U_{n-1}(x).$$

Define  $\mathcal{D}_q^{-1}$  by

$$\mathcal{D}_q^{-1} \sum_{n=0}^{\infty} f_n U_n(x) = \sum_{n=1}^{\infty} f_{n-1} \frac{q^{1/2} - q^{-1/2}}{q^{n/2} - q^{-n/2}} T_n(x).$$

This defines a map from  $L_2[(1 - x^2)^{1/2}, -1, 1]$  into  $L_2[(1 - x^2)^{-1/2}, -1, 1]$ . Indeed it has the representation

$$\begin{aligned} & (\mathcal{D}_q^{-1} g)(\cos \theta) \\ &= \frac{1 - q}{4\pi\sqrt{q}} \int_{-\pi}^{\pi} \frac{\vartheta'_4((\theta + \phi)/2, \sqrt{q})}{\vartheta_4((\theta + \phi)/2, \sqrt{q})} g(\cos \theta) \sin \theta d\theta. \end{aligned}$$

Moreover  $\mathcal{D}_q \mathcal{D}_q^{-1} = \text{identity}$ .

One can diagonalize this integral operator. Similarly for  $q$ -ultraspherical polynomials, Ismail–Zhang, and the continuous  $q$ -Jacobi polynomials, Ismail–Rahman–Zhang. The kernel for the full Askey–Wilson is also known, Ismail–Rahman. This led to a  $q$ -analogue of the exponential function.

## Taylor Series.

The Askey–Wilson basis is  $\{(ae^{i\theta}, ae^{-i\theta}; q)_n\}$ .

Also

$$\begin{aligned} & \mathcal{D}_q(ae^{i\theta}, ae^{-i\theta}; q)_n \\ &= -\frac{2a(1-q^n)}{1-q} (aq^{1/2}e^{i\theta}, aq^{1/2}e^{-i\theta}; q)_{n-1}. \end{aligned}$$

**Theorem 2.** *Let  $f$  be a polynomial, then*

$$f(x) = \sum_{k=0}^n f_k(ae^{i\theta}, ae^{-i\theta}; q)_k,$$

where

$$f_k = \frac{(q-1)^k}{(2a)^k (q; q)_k} q^{-k(k-1)/4} (\mathcal{D}_q^k f)(x_k)$$

with

$$x_k := \frac{1}{2}(aq^{k/2} + q^{-k/2}/a).$$

When  $f(x) = (be^{i\theta}, be^{-i\theta}; q)_n$  we get the  $q$ -Pfaff–Saalschütz theorem.

Can we build the theory of  $q$ -series this way?

Extension to non polynomial cases. This is an interpolation problem.

A classic is to reconstruct an entire function from its values at the integers. Not always possible because  $\frac{\sin(\pi z)}{\pi z} = 0$  at  $z = \pm 1, \pm 2, \dots$ .

The entire function is unique if  $\log M(f, r) = cr[1 + o(1)]$ ,  $c < \pi$ . Gelfond and his school considered interpolation at  $q^{-n}$ ,  $|q| < 1$ . Our interpolation points are

$$x_k := \frac{1}{2}(aq^{k/2} + q^{-k/2}/a).$$

There is an expansion theorem when  $M(f, r)$  has certain growth condition (Ismail–Stanton).

S. Cooper 1996 proved by induction

$$\mathcal{D}_q^n f(x) = \frac{(2z)^n q^{n(3-n)/4}}{(q-1)^n} \times \sum_{k=0}^n \begin{bmatrix} n \\ k \end{bmatrix}_q \frac{q^{k(n-k)} z^{-2k} \bar{f}(q^{k-n/2} z)}{\left( q^{n-2k+1} z^{-2}; q \right)_k \left( z^2 q^{2k+1-n}; q \right)_{n-k}}.$$

Application 1: Apply this to

$$f(\cos \theta) = \frac{(\alpha e^{i\theta}, \alpha e^{-i\theta}; q)_\infty}{(\beta e^{i\theta}, \beta e^{-i\theta}; q)_\infty}$$

you get the  ${}_6\phi_5$  summation theorem. The Rodrigues formula is

$$\begin{aligned} & w(x; \mathbf{a} | q) p_n(x; \mathbf{a} | q) \\ &= \left( \frac{q-1}{2} \right)^n q^{n(n-1)/4} \mathcal{D}_q^n \left[ w(x; q^{n/2} \mathbf{a} | q) \right]. \end{aligned}$$

Application 2: The Rodrigues formula for the Askey–Wilson polynomials gives the  ${}_8\phi_7$  to  ${}_4\phi_3$  transformation (Watson).

Expansions in the Askey–Wilson polynomials.

$$f(x) = \sum_{n=0}^{\infty} f_n p_n(x; \mathbf{a}).$$

If  $\mathcal{D}_q f$  acts nicely on  $f$  we can find the coefficients using the Rodrigues formula.

**Theorem 3.** (Ismail–Stanton 2013) *We have the following expansion*

$$\begin{aligned}
& {}_{p+1}\phi_p \left( \begin{matrix} a_1, \dots, a_{p-1}, t_4 e^{i\theta}, t_4 e^{-i\theta} \\ t_1 t_4, t_2 t_4, t_3 t_4, b_1, \dots, b_{p-3} \end{matrix} \middle| q, \zeta \right) \\
&= \sum_{k=0}^{\infty} p_k(\cos \theta; \mathbf{t} | q) \frac{(a_1, \dots, a_{p-1}; q)_k}{(t_1 t_4, t_2 t_4, t_3 t_4, b_1, \dots, b_{p-3}; q)_k} \\
&\quad \times \frac{(-t_4 \zeta)^k q^{\binom{k}{2}}}{(q, t_1 t_2 t_3 t_4 q^{k-1}; q)_k} \\
&\quad \times {}_{p-1}\phi_{p-2} \left( \begin{matrix} q^k a_1, \dots, q^k a_{p-1} \\ q^k b_1, \dots, q^k b_{p-3}, t_1 t_2 t_3 t_4 q^{2k} \end{matrix} \middle| q, \zeta \right).
\end{aligned}$$

The Andrews formula (2012) is the case  $p = 4$  in Theorem 3 with the parameter identification

$$a_1 = q^{-N}, a_2 = \rho_1, a_3 = \rho_2, b_1 = \rho_1 \rho_2 q^{-N} / a, \zeta = q.$$

In this case the  ${}_3\phi_2$  can be summed by the  $q$ -Pfaff–Saalschütz theorem.

Fields and Wimp (1961).

$$\begin{aligned}
 & \sum_{m=0}^{\infty} a_m b_m \frac{(zw)^m}{m!} \\
 = & \sum_{n=0}^{\infty} \frac{(-w)^n}{n! (\gamma + n)_n} \left( \sum_{r=0}^{\infty} \frac{b_{n+r} w^r}{r! (\gamma + 2n + 1)_r} \right) \\
 & \times \left[ \sum_{s=0}^n \frac{(-n)_s (n + \gamma)_s}{s!} a_s z^s \right].
 \end{aligned}$$

This version is due to Verma 1972. Lagrange inversion (Gessel–Stanton).

The following general expansion follows from  $q$ -Dixon's theorem ( $4\phi_3$ )

$$\begin{aligned}
 & \sum_{n=0}^{\infty} A_n B_n \frac{(t_4 z, t_4/z; q)_n}{(q; q)_n} \zeta^n \\
 = & \sum_{k=0}^{\infty} \frac{(-\zeta)^k q^{\binom{k}{2}}}{(q, Cq^{k-1}; q)_k} \left[ \sum_{n=0}^{\infty} \frac{B_{n+k} \zeta^n}{(q, Cq^{2k}; q)_n} \right] \\
 \times & \left[ \sum_{j=0}^k \frac{(q^{-k}, Cq^{k-1}; q)_j}{(q; q)_j} A_j (t_4 z, t_4/z; q)_j q^j \right].
 \end{aligned}$$



Moments of the Askey–Wilson weight functions. Used by Corteel–Stanley–Stanton–Williams.

The Stieltjes electrostatic equilibrium leads to system is

$$\frac{\beta + 1}{1 + x_j} - \frac{\alpha + 1}{1 - x_j} + \sum_{1 \leq k \leq n, k \neq j} \frac{2}{x_j - x_k} = 0,$$

for  $1 \leq j \leq n$ . Stieltjes used  $y = \prod_{j=1}^n (x - x_j)$  and

$$\sum_{1 \leq k \leq n, k \neq j} \frac{2}{x_j - x_k} = \frac{y''(x_j)}{y'(x_j)}.$$

He turned this system to

$$(1 - x^2)y'' + (\beta - \alpha - x(\alpha + \beta + 2))y' + \lambda_n y = 0,$$

at  $x = x_1, \dots, x_n$ . By choosing  $\lambda$  this becomes valid for all  $x$ . Thus the equilibrium points are at the zeros of a Jacobi polynomial  $P_n^{(\alpha, \beta)}(x)$ .

For the XXZ model the Bethe Ansatz equations are ( $1 \leq k \leq n$ ),

$$= \prod_{j \neq k, j=1}^n \frac{\left( \frac{\sin \left( \lambda_k + \frac{1}{2}\eta \right)}{\sin \left( \lambda_k - \frac{1}{2}\eta \right)} \right)^{2N} \sin \left( \lambda_k + \lambda_j + \eta \right) \sin \left( \lambda_k - \lambda_j + \eta \right)}{\sin \left( \lambda_k + \lambda_j - \eta \right) \sin \left( \lambda_k - \lambda_j - \eta \right)}.$$

Change the system of equations to

$$= \prod_{j \neq k, j=1}^n \frac{\prod_{\ell=1}^{2N} \frac{\sin \left( \lambda_k + s_\ell \eta \right)}{\sin \left( \lambda_k - s_\ell \eta \right)} \sin \left( \lambda_k + \lambda_j + \eta \right) \sin \left( \lambda_k - \lambda_j + \eta \right)}{\sin \left( \lambda_k + \lambda_j - \eta \right) \sin \left( \lambda_k - \lambda_j - \eta \right)},$$

for  $1 \leq k \leq n$ , where  $s_\ell$ 's are  $2N$  complex numbers. Notation:

$$q = e^{2i\eta}, \quad \theta = 2\lambda, \quad t_j = q^{-s_j}.$$

$$w(\cos \theta; t_1, \dots, t_{2N}) \\ := \frac{(e^{iN\theta}, e^{-iN\theta}; q^{N/2})_\infty}{\sin(N\theta/2) \prod_{j=1}^{2N} (t_j e^{i\theta}, t_j e^{-i\theta}; q)_\infty}.$$

**Notation**  $\mathbf{t} = (t_1, t_2, \dots, t_{2N})$ .

$$\frac{1}{w(x; \mathbf{t})} \mathcal{D}_q \left( w(x; q^{1/2} \mathbf{t}) \mathcal{D}_q \right) y(x) = r(x) y(x),$$

where  $r(x)$  is a polynomial of degree  $\leq N - 2$ .

$$\Pi(z; \mathbf{t}) \mathcal{D}_q^2 y + \Phi(z; \mathbf{t}) \mathcal{A}_q \mathcal{D}_q y = r(x) y.$$

The solution of the Bethe Ansatz equations is at the zeros of the polynomial solutions.

Joint work with Lin and Roan.